### The missing 95%: Theory and Phenomenology of Dark Matter and Dark Energy

Joachim Kopp

DPG Spring Meeting Göttingen, March 2012

# **‡Fermilab**

### Outline



#### 2 Finding dark matter

- Direct detection
- Indirect detection
- Production at colliders

#### 3 Modelling dark matter

#### 4 Dark energy

### Outline



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#### Dark energy

### **Celestial mechanics**

Stellar/galactic dynamics relates:

- The mass distribution (inferred from brightness)
- Kinetic energy (inferred from Doppler shifts)



Fritz Zwicky 1898–1974



Vera Rubin 1928–

#### Observations of rotational velocities in galaxies show: Rubin 1975



The gravitational pull on peripheral stars is stronger than predicted from the mass of the luminous matter *M* 

$$m\frac{v^2}{r}=G_N\frac{mM}{r^2}$$

at  $r \to \infty$ 

### Collisions of galaxy clusters



Artist's rendering (Image: NASA)

red = gas (from x-ray observations) blue = (dark) matter distribution (from gravitational lensing)

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Image: NASA (Chandra [x-ray], ESO WFI [lensing], HST [optical])

red = gas (from x-ray observations) blue = (dark) matter distribution (from gravitational lensing)

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WMAP's observation of the CMB: A fingerprint of the universe at  $t \simeq 300\,000$  yrs

(when electrons and protons first combined to form atoms).





red = overdense, hot regions  $(0 \dots + 200 \ \mu K)$ blue = underdense, cold regions  $(-200 \dots 0 \ \mu K)$ 

Image credit: NASA

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#### red curve = theory prediction black points = WMAP data

Image credit: NASA

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Image credit: NASA



## What is this stuff?

- Modified laws of gravity?
  - Hard to explain all observations
- MACHOs (Massive Compact Halo Objects)?
  - Planets, Brown dwarfs, neutron stars, ...
  - ▶ Ruled out as dark matter in the mass range  $0.6 \times 10^{-7} M_{\odot} < M < 15 M_{\odot}$  by searches for gravitational microlensing
  - Searches for candidate objects yield too few of them
- Hot (relativistic) Dark Matter (neutrinos or other relativistic particles)?
  - Cannot explain large scale structure of the universe (hot dark matter would smoothen the galaxy distribution)

#### Cold or Warm Dark Matter

- Axions
  - Ultra-light, but non-relativistic due to non-thermal production
- Gravitinos
  - \* Only gravitational couplings  $\rightarrow$  bad for direct/indirect/collider detection
- WIMPs (Weakly Interacting Massive Particle)
  - \* New, heavy, stable particles
  - Should have some non-gravitational interaction with SM particles for production in the early universe

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#### **Direct Dark Matter detection**

Idea: A WIMP (Weakly Interacting Massive Particle) can scatter on an atomic nucleus.



Strategy: Look for feeble nuclear recoil

Problem: Many background processes (radioactive decays, cosmic rays, ...) can mimic the signal





#### **Direct detection results**



Assumptions here: Elastic DM scattering  $\propto$  target mass (often realized in SUSY)

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- Previous slide: Elastic dark matter ( $\chi$ ) scattering through scalar current  $[(\bar{q}q)(\bar{\chi}\chi)]$  or vector current  $[(\bar{q}\gamma_{\mu}q)(\bar{\chi}\gamma^{\mu}\chi)]$  assumed  $\Rightarrow$  Cross section  $\propto$  target mass
- In models with different coupling structure, the relative detection efficiencies of different experimental technologies may be different

- Spin-dependent couplings
  - E.g. coupling through axial vector current  $[(\bar{q}\gamma^{\mu}\gamma^{5}q)(\bar{\chi}\gamma_{\mu}\gamma^{5}\chi)]$
  - Cross section 
    <u>x</u> target spin
  - Cannot explain DAMA, CoGeNT, CRESST results

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- Inelastic dark matter Tucker-Smith Weiner hep-ph/0101138
  - There may be two DM states  $\chi$  and  $\chi'$  with  $m'_{\chi} = m_{\chi} + \delta$  ( $\delta \sim 100 \text{ keV}$ )
  - ► Scattering  $\chi N \rightarrow \chi' N \Rightarrow$  heavy target nuclei kinematically preferred
  - Could explain CRESST, but not DAMA JK Schwetz Zupan 1110.2721



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• . . .

# Conclusion: Hard to explain all data simultaneously

### Direct detection uncertainties

• Large uncertainty in local DM density

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- Large uncertainties in DM velocity distribution
  - Scattering rate depends strongly on DM velocity
  - DM streams?
  - Debris flow?



Kuhlen Lisanti Spergel arXiv:1202.0007, graphics courtesy of Mariangela Lisanti

#### Direct detection uncertainties

- Large uncertainty in local DM density
- Large uncertainties in DM velocity distribution
  - Scattering rate depends strongly on DM velocity
  - DM streams?
  - Debris flow?
- Predicting WIMP–nucleus cross sections is difficult
  - Models predict WIMP-quark cross section
  - Need to know quark content of the nucleon
  - ► Especially problematic for Higgs-mediated scattering: coupling ∝ quark mass ⇒ sea quarks dominate
  - Need to know nuclear form factor especially difficult for spin-dependent scattering

### Outline

Evidence for dark matter



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3 Modelling dark matter

#### Dark energy

### Indirect Dark Matter detection

Idea: WIMPs (Weakly Interacting Massive Particles)  $\chi$  can annihilate (or decay) into Standard Model particles (*f*) in an astrophysical environment.



Strategy: Look for annihilation products in cosmic rays

Problems:

- Many other sources of cosmic rays
- Propagation of charged particles in the galaxy poorly understood

Advantage:

Many sources to look at









look at many sources

















### Indirect DM detection — Examples

#### $\gamma$ -rays from dwarf galaxies



#### Idea:

Look for anomalous  $\gamma$ -ray flux

Pro:

Few stars  $\Rightarrow$  few backgrounds

#### Con:

- Relatively low DM density
- Results model-dependent
- Large astrophysical uncertainties
# Indirect DM detection — Examples

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# Indirect DM detection — Examples

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#### Other indirect DM searches:

- Cosmic anti-matter (e<sup>+</sup>,  $\bar{p}$ , ...) PAMELA, Fermil-LAT, ...
- γ-rays from the galactic center Hooper et al.
- High-energy neutrinos from the Sun IceCube, SuperKamiokande

ο..

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## Dark matter at colliders



### Dark matter at colliders



make your own needles!





## Generic collider searches for dark matter

Idea:

• Produce WIMPs in collisions of Standard Model particles



WIMPs can recoil against a jet or a photon from initial state radiation



■ Experimental signatures: Mono-jets + ∉<sub>T</sub> and mono-photons + ∉

## LHC limits on DM-quark couplings



- Assumptions here:
  - Effective field theory approach valid (limits may be better or worse if EFT not valid)
  - Equal coupling to all quark flavors
- Extremely competitive limits for
  - Light dark matter (below direct detection threshold)
  - DM coupled to gluons (high gluon luminosity at the LHC)
  - Spin-dependent DM interactions (DD suffers from loss of coherence)

## Model-dependent collider searches: SUSY-DM

Idea:

- In many models, DM is produced in the decay of heavy, strongly interacting particles (for instance squarks and gluinos in SUSY)
- Experimental signature: something + missing energy
- Example:  $pp \rightarrow (\tilde{g} \rightarrow jZ\chi^0)(\tilde{q} \rightarrow jjW\chi^0)$

- Advantage: Very sensitive
- Problem:

Minor modifications to the model may drastically change the phenomenology

 Problem (all collider searches): Collider can only find DM candidate(s)



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- Conclusion: If dark matter originates from electroweak-scale new physics, it automatically has the right abundance

# The Wimp Miracle



 Motivation: The DM and baryon energy densities in the universe are similar

 $\Omega_{DM}\simeq 5\,\Omega_b$ 

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- Baryon density Ω<sub>b</sub> generated by yet unknown dynamics behind the particle–antiparticle asymmetry of the universe (not by thermal freeze-out)
- Assume dark matter ( $\chi$ ) density is also determined by  $\bar{\chi}$ - $\chi$  asymmetry  $\Rightarrow$  Asymmetric dark matter

# Models of asymmetric dark matter

Example 1 Kaplan Luty Zurek, arXiv:0901.4117

- B L asymmetry generated at high T (e.g. via Leptogenesis)
- Effective superfield operator

$$\mathcal{L} \supset \frac{1}{M} \bar{X}^2 L H_u$$
 (\*)

transfers  $B - L \leftrightarrow 2X$ , e.g. via



• Final X (DM number) asymmetry depends on # of SM species contributing to (\*) at freeze-out

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#### Example 2

Buckley Randall 1009.0270 Blennow et al. 1009.3159

- Generate X asymmetry in hidden sector
- Transfer to *B L* asymmetry in the SM sector
  - via B L violating interactions (e.g. (\*))
    - via sphaleron processes

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#### Example 3

Davoudiasl et al. 1008.2399 Gu Lindner Sarkar Zhang 1009.2690

- New heavy particles decay partly into DM, partly into SM particles
- B L X is conserved
- DM (X) does not participate in SM sphaleron processes
   ⇒ Asymmetry frozen in

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# Evidence for dark energy: Type Ia Supernovae

- When a white dwarf accretes matter from a companion star, it becomes unstable once it reaches  $\sim 1.4 M_{\odot}$ 
  - Re-ignition of nuclear fusion
  - Thermonuclear explosion
- Since the progenitor mass is always ~ 1.4M<sub>☉</sub>, all Type Ia Supernovae are very similar
  - Energy release precisely known
  - SN la are standard candles
- Measurement:
  - ► Apparent brightness → distance
  - ▶ Redshift → velocity
- Result:
  - Long ago (very distant SN Ia, low brightness), the universe was expanding more slowly than we thought!
  - It must be accelerating
- CMB and Large Scale Structure observations confirm this





## What is accelerating the Universe?

- A cosmological constant?
  - An ad-hoc addition to the Einstein equations

$$R_{\mu\nu}-\frac{1}{2}g_{\mu\nu}R_{\alpha}^{\ \alpha}=8\pi G\,T_{\mu\nu}+g_{\mu\nu}\Lambda$$

- Observations require Λ ~ (10<sup>-12</sup> GeV)<sup>4</sup>
- Extra source of energy with negative pressure

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- Observations require Λ ~ (10<sup>-12</sup> GeV)<sup>4</sup>
- Extra source of energy with negative pressure
- QFT vacuum energy?
  - A vacuum expectation value (vev) or condensate of a quantum field behaves like a cosmological constant
  - ► Problem: All known condensates/vevs are way too large! (We expect  $\Lambda \sim M_{Pl}^4 \sim (10^{19} \text{ GeV})^4$ )

## What is accelerating the Universe? (cont'd)

- Quintessence: A new, slowly rolling scalar field
  - Introduce new scalar field  $\phi$  slowly rolling down its potential  $V(\phi)$
  - Lagrangian:

$$\mathcal{L}_{\phi} = rac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi)$$

Energy and pressure:

$$p = \frac{1}{2}\dot{\phi}^2 + V(\phi),$$
  $p = \frac{1}{2}\dot{\phi}^2 - V(\phi)$ 

• A cosmological constant corresponds to  $\rho = -p \Rightarrow$  require  $\dot{\phi}^2 \ll V(\phi)$ 



for a review see Caldwell Kamionkowski 0903.0866

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• A cosmological constant corresponds to  $\rho = -p \Rightarrow$  require  $\dot{\phi}^2 \ll V(\phi)$ • Extensions of general relativity

Scalar-tensor gravity: Modified Einstein-Hilbert action

$$S = \frac{1}{16\pi G} \int \sqrt{-g} \, d^4 x \, R \quad \rightarrow \quad S = \frac{1}{16\pi G} \int \sqrt{-g} \, d^4 x \, f(\phi) \times R$$

A special case: f(R) gravity:

$$S=\frac{1}{16\pi G}\int\sqrt{-g}\,d^4x\,f(R)$$

## Summary

- Overwhelming evidence for dark matter
- A lot of data available
  - Direct detection
    - \* Difficult to reconcile possible evidence with null results
  - Indirect searches
    - Strong exclusion limits
    - \* Suffers from poorly understood astrophysical backgrounds
  - Collider searches
    - ★ Generic searches (monojets + ∉<sub>T</sub>, mono-γ + ∉) and model-specific searches (cascade decays) are underway full-steam
- Dark matter models
  - Dark matter from electroweak scale new physics: Correct cosmic abundance due to WIMP Miracle
  - Light (10 GeV) dark matter: Correct cosmic abundance if related to baryon-antibaryon asymmetry
- Dark energy
  - Accelerated expansion of the Universe well-established
  - So far, a cosmological constant is the leading explanation

Thank you!

### **Bonus** material

## Spin-dependent DM couplings?

- Previous slide: Dark matter  $(\chi)$  couplings through scalar current  $[(\bar{q}q)(\bar{\chi}\chi)]$  or vector current  $[(\bar{q}\gamma_{\mu}q)(\bar{\chi}\gamma^{\mu}\chi)]$  assumed  $\Rightarrow$  Cross section  $\propto$  target mass
- Alternative: Axial vector  $[(\bar{q}\gamma^{\mu}\gamma^{5}q)(\bar{\chi}\gamma_{\mu}\gamma^{5}\chi)]$  interaction
  - $\Rightarrow$  Cross section  $\propto$  target spin



Note: CoGeNT & CRESST have very low sensitivity to spin-dependent DM scattering.

## Inelastic dark matter?

Idea: There may be two DM states  $\chi$  and  $\chi'$  with

 $m_{\chi'} = m_{\chi} + \delta$ 

Scattering proceeds via

 $\chi + \mathbf{N} \rightarrow \chi' + \mathbf{N}$ 

- Modified kinematics compared to elastic scattering
- Affects different target nuclei differently

## Inelastic dark matter?

#### Idea: There may be two DM states $\chi$ and $\chi'$ with

 $m_{\chi'} = m_{\chi} + \delta$ 



plot from JK Schwetz Zupan 1110.2721
## Isospin-violating dark matter?

Idea: Dark matter could couple differently to protons and neutrons  $\Rightarrow$  Detection efficiencies of different target materials change



plot from JK Schwetz Zupan 1110.2721

 $f_n$ ,  $f_p$ : DM couplings to protons and neutrons  $A_{\text{eff}}$ : Effective nuclear mass for DM scattering

## Isospin-violating dark matter?

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plot from JK Schwetz Zupan 1110.2721

Idea: DM could couple only to leptons at tree level

- DAMA and CoGeNT do not reject electron-recoils as background
- But: Electron recoils above threshold (≥ 1 keV) strongly suppressed (electron needs large initial momentum → probe high-p tail of wave functions)



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- Thus: DM-nucleus scattering dominates, even if loop-induced



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- $\bullet\,$  But: Electron recoils above threshold ( $\gtrsim$  1 keV) strongly suppressed
- Thus: DM-nucleus scattering dominates, even if loop-induced
- But: Loop diagrams forbidden for some models
- Problems then:
  - Very large couplings needed to compensate wave function suppression
  - Poor fit to DAMA and CoGeNT energy spectra



## The Galactic Center



#### Pros:

Highest DM density

Cons:

- DM distribution uncertain
- Many background sources

## The Galactic Center



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## **Dwarf Galaxies**



### Pros:

Few backgrounds

#### Cons:

Relatively low DM density

## The Galactic Center



### Pros:

Highest DM density

### Cons:

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- Many background sources

## **Dwarf Galaxies**



## Cosmic antimatter



### Pros:

Few background sources

### Cons:

- Backgrounds uncertain
- Propagation of charged particle has large uncertainties
- Non-directional



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- Backgrounds uncertain
- Propagation of charged particle has large uncertainties
- Non-directional

# High-energy neutrinos



## Idea:

- DM capture/annihilation in the Sun
- Flux dominated by capture rate
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Low neutrino cross sections



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IceCube collaboration, 1111.2738

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# What is a sphaleron?

 SU(2) gauge field vacuum configurations are classified according to their winding number (or Chern-Simons number)

$$N_{CS} = \frac{1}{16\pi^2} \int_0^t dt \int d^3 x \operatorname{tr} F_{\mu\nu} \tilde{F}^{\mu\nu} \qquad \underbrace{\overbrace{ind_{r}(z_0) = -1}^{z_0}}_{\operatorname{Ind_{r}(z_0) = +1}} \underbrace{\overbrace{ind_{r}(z_0) = 0}^{z_0}}_{\operatorname{Ind_{r}(z_0) = +2}}$$

Configurations with different winding number cannot be continuously transformed into each other.

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- Sphalerons are processes (with *E* > 0) that change the winding number Their energy is of order *m<sub>H</sub>*, the symmetry breaking scale (100 GeV)

## What is a sphaleron?

- *SU*(2) gauge field vacuum configurations are classified according to their winding number (or Chern-Simons number)  $N_{CS} = \frac{1}{16\pi^2} \int_0^t dt \int d^3x \operatorname{tr} F_{\mu\nu} \tilde{F}^{\mu\nu}$
- Sphalerons are processes (with E > 0) that change the winding number
- In the SM, a change in winding number corresponds to a change in B+L. In fact, considering only left-handed (SU(2)<sub>L</sub>-charged) fermions:

$$j^{\mu}_{B+L} = \sum_{\psi=q,\ell} \frac{1}{2} \bar{\psi} \gamma^{\mu} (1-\gamma^5) \psi$$

A change in B + L is equivalent to a change in  $N_{CS}$ :