# Density Functional Theory: from theory to Applications

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Fully atomistic QM/MM approaches Embedding in continuum environment An example of QM/MM calculations

#### Fully atomistic QM/MM approaches CP-PAW Interface CPMD/ GROMOS Interface Real space multi-grid approach

#### Embedding in continuum environment CP-PAW/COSMO Interface

An example of QM/MM calculations

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# Mixed quantum/ classical molecular dynamics (QM/MM)



- Quantum mechanics is computationally expensive
- Idea: subdivide a complex system in a small, relevant part (QM) and the environment which is not neglected but treated at lower level of accuracy (MM).<sup>1</sup>

<sup>1</sup>See pioneering work of Warshel and Levitt, JMB 1976, 103, 227. ( = ) = 🔊 a 🖉

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Metallo-beta-lactamase enzyme. Two zinc ions (purple spheres) are in the enzyme's active site (with amino acids coordinating the metals represented as sticks).

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Fully atomistic QM/MM approaches: both QM and MM are based on particle representation. The full Hamiltonian is:

$$E^{QM/MM} = E^{QM}(\{\mathbf{R}_{\alpha}\}) + E^{MM}(\{\mathbf{R}_{I}\}) + E^{QM-MM}(\{\mathbf{R}_{\alpha}\}, \{\mathbf{R}_{I}\})$$
(1)

- ► E<sup>QM</sup>({**R**<sub>α</sub>}) is the Kohn-Sham Hamiltonian H<sup>KS</sup><sub>e</sub>, e.g. in a plane wave / pseudopotential representation.
- ► E<sup>MM</sup>({**R**<sub>I</sub>}) are force field energy expressions: short-range bonded interaction with fixed topology + non-bonded interaction (electrostic and van der Waals)

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#### MM: Bonds and bends



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#### MM: Urey-Bradley and Dihedral



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### MM: Improper (Out of plane motion)





By definition, from the force field parameter and documentation, the first atom listed is the central atom. Therefore just like a normal torsion, the angle of interest is between the plane defined of particle *ijk* to that of the plane defined by particle *jkl*. This is only used for special situations.

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#### MM: Non-bonded

 Steric interactions are usually described by standard Lenard-Jones interaction, but different functional forms are also possible



Electrostatics:

$$V(r_{ij}) = \frac{q_i q_j}{\epsilon r_{ij}} \tag{2}$$

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### Biomolecular simulations: AMBER

$$\begin{split} V_{total} &= \sum_{bonds} K_b (b - b_{eq})^2 + \\ &\sum_{angle} K_\theta (\theta - \theta_{eq})^2 + \\ &\sum_{dihedrals} \frac{V_n}{2} \big[ 1 + \cos(n\phi - \gamma) \big] + \\ &\sum_{i < j} \left[ \frac{A_{ij}}{R_{ij}^{12}} - \frac{B_{ij}}{R_{ij}^6} - \frac{q_i q_j}{\varepsilon R_{ij}} \right] \end{split}$$



#### http://ambermd.org/

D.A. Case, T.E. Cheatham, III, T. Darden, H. Gohlke, R. Luo, K.M. Merz, Jr., A. Onufriev, C. Simmerling, B. Wang and R. Woods. The Amber biomolecular simulation programs. J. Computat. Chem. 26, 1668-1688 (2005).

J.W. Ponder and D.A. Case. Force fields for protein simulations. Adv. Prot. Chem. 66, 27-85 (2003).

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#### Biomolecular simulations: CHARMM

#### http://www.charmm.org

CHARMM: A Program for Macromolecular Energy, Minimization, and Dynamics Calculations, J. Comp. Chem. 4, 187-217 (1983), by B. R. Brooks, R. E. Bruccoleri, B. D. Olafson, D. J. States, S. Swaminathan, and M. Karplus.

CHARMM: The Energy Function and Its Parameterization with an Overview of the Program, in The Encyclopedia of Computational Chemistry, 1, 271-277, P. v. R. Schleyer et al., editors (John Wiley & Sons: Chichester, 1998), by A. D. MacKerell, Jr., B. Brooks, C. L. Brooks, III, L. Nilsson, B. Roux, Y. Won, and M. Karplus,

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# OPLS (Optimized Potentials for Liquid Simulations)

$$\begin{split} V_{total} &= \sum_{bonds} K_r (r - r_{eq})^2 + \\ &\sum_{angle} K_{\theta} (\theta - \theta_{eq})^2 + \\ &\sum_{dihedrals} \frac{V_1}{2} [1 + \cos(\phi - \gamma_1)] + \\ &\frac{V_2}{2} [1 + \cos(2\phi - \gamma_2)] + \\ &\frac{V_3}{2} [1 + \cos(3\phi - \gamma_3)] + \\ &\sum_{l < j} \left[ 4\varepsilon_{ij} \left( \frac{\sigma_{ij}^{12}}{r_{ij}^{12}} - \frac{\sigma_{ij}^6}{r_{ij}^6} \right) - \frac{q_i q_j e^2}{r_{ij}} \right] f_{ij} \end{split}$$

Jorgensen WL, Tirado-Rives J (1988). "The OPLS Force Field for Proteins. Energy Minimizations for Crystals of Cyclic Peptides and Crambin". J. Am. Chem. Soc. 110: 16571666 Jorgensen WL, Maxwell DS, Tirado-Rives J (1996). "Development and Testing of the **OPLS All-Atom Force Field on** Conformational Energetics and Properties of Organic Liquids". J. Am. Chem. Soc. 118 (45): 1122511236

 $\begin{array}{l} \gamma_1 = \gamma_2 = \gamma_3 = 0^{\circ} \\ \varepsilon_{ij} = \sqrt{\varepsilon_{ii}\varepsilon_{jj}} \end{array} \qquad \sigma_{ij} = \sqrt{\sigma_{ii}\sigma_{jj}} \quad f_{ij} = 0.5 \text{ for } 1 - 4 \text{ interactions} \\ f_{ij} = 1.0 \text{ for everything else} \end{array}$ 

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#### Class II force fields

- Greater transferability: Application of empirical force field parameters to molecules not explicitly included during the parameter optimization.
- Example: Merck Molecular Force Field (MMFF), a Class II force field designed to be a transferable force field for pharmaceutical compounds that accurately treats conformational energetics and non-bonded interactions. This would, ideally, produce a force field that was adequate for both gas phase and condensed phase calculations.

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#### Class II force fields: MMFF

$$\begin{split} V_{total} &= \sum_{bonds} K_{bond} \left( r - r_{eq} \right)^2 \left( 1 + cs \left( r - r_{eq} \right) + \frac{7}{12} \left( cs^2 \left( r - r_{eq} \right)^2 \right) \right) + \\ &\sum_{angle} K_{\theta} \left( \theta - \theta_{eq} \right)^2 \left( 1 + cb \left( \theta - \theta_{eq} \right) \right) + \\ &\sum_{angle, linear} K_{al} \left( 1 + \cos(\theta) \right) + \\ &\sum_{angle, linear} \left( K_{ijk} \left( r_{ij} - r_{eq} \right) + K_{kji} \left( r_{kj} - r_{eq} \right) \right) \left( \theta - \theta_{eq} \right) + \\ &\sum_{stretch, bend} K_{OOP} (\chi)^2 + \\ &\sum_{dihedrals} \frac{V_1}{2} \left[ 1 + \cos(\phi) \right] + \frac{V_2}{2} \left[ 1 + \cos(2\phi) \right] + \frac{V_3}{2} \left[ 1 + \cos(3\phi) \right] + \\ &\sum_{i < j} \left[ \varepsilon_{ij} \left( \frac{1.07\sigma}{r_{ij} + 0.07\sigma} \right)^7 \left( \frac{1.12\sigma^7}{r_{ij}^7 + 0.07\sigma^7} - 2 \right) - \frac{q_i q_j}{D(r_{ij} + \delta)} \right] \end{split}$$

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# $E^{QM-MM}(\{\mathbf{R}_{\alpha}\},\{\mathbf{R}_{I}\})$

$$E^{QM-MM} = E_b^{QM-MM} + E_{nb}^{QM-MM} \tag{3}$$

$$E_{nb}^{QM-MM} = E_{es}^{QM-MM} + E_{steric}^{QM-MM}$$
(4)

 E<sup>QM-MM</sup><sub>steric</sub> generally follows the model used in E<sup>MM</sup> (Lenard-Jones-type potentials).

• For the  $E_{es}^{QM-MM}$  there different coupling schemes:

- mechanical embedding: no influence of MM charges on the QM system
- electrostatic embedding: electrostatic interaction between mm charges and charge density of the QM system
- polarized embedding: MM polarization due to the QM system is also included

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### Technical issues connected with the electrostic coupling

Straightforward numerical evaluation of the electrostatic coupling term:

$$E_{es}^{QM-MM} = \sum_{I \in MM} q_I \int \frac{n(r)}{|\mathbf{r} - \mathbf{R}_I|} d\mathbf{r}$$
(5)

is prohibitive. Indeed it would involve a number of operations which scales as  $N_g$  (number of grid points) times  $N_{MM}$  (number of MM atoms).

Electron spill out or charge leakage

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### **CP-PAW** Interface

Pioneering qm/mm work by Bloch and coworkers<sup>2</sup>. Motivation: studying chemical reaction of complex system and computing free energy profile using thermodynamics integration. The instaneous electronic density of the QM system is mapped onto a set of Gaussian smearing functions tied to the QM sites. The corresponding Lagrangian is:

$$\mathcal{L}^{QM/MM} = \mathcal{L}^{QM}_{CP} + \mathcal{L}^{MM} + \mathcal{L}^{QM-MM}$$
(6)  
$$= \sum_{\alpha} \frac{1}{2} M_{\alpha} \dot{\mathbf{R}}^{2}_{\alpha} + \sum_{i} \mu \left\langle \dot{\phi}_{i} | \dot{\phi}_{i} \right\rangle - \left\langle \Psi_{0} | \tilde{H}^{KS}_{e}(\{\mathbf{R}_{\alpha}\}) | \Psi_{0} \right\rangle (7)$$
$$+ \sum_{i,j} \Lambda_{i,j}(\langle \phi_{i} | \phi_{j} \rangle - \delta_{ij})$$
(8)  
$$+ \sum_{l} \frac{1}{2} M_{l} \dot{\mathbf{R}}^{2}_{l} + E^{MM}(\{\mathbf{R}_{l}\}) + E^{QM-MM}(\{\mathbf{R}_{l}\}, \{\mathbf{R}_{\alpha}\})$$
(9)

<sup>2</sup>J. Chem. Phys. 103, 7422

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### **CP-PAW** Interface

- Separate Nose'-Hoover thermostats for the QM/MM regions
- Decoupling scheme between periodic images
- The continuous charge density is compressed to a linear superposition of of atom-centered spherical Gaussian functions:

$$n^{G}(\mathbf{r}) = \sum_{\alpha} q_{\alpha} G(\mathbf{r} - \mathbf{R}_{\alpha})$$
(10)

Gaussian parameters are determined imposing that  $n^{G}(\mathbf{r})$  reproduce the multiple moments of  $n(\mathbf{r})$ .

In the limit of infinitesimally small width the Gaussian density becomes:

$$n^{pc}(\mathbf{r}) = \sum_{\alpha} q_{\alpha} \delta(\mathbf{r} - \mathbf{R}_{\alpha})$$
(11)

The electrostatic coupling is obtained between these effective QM charges and the MM charges.

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# Handling bond cuts: the hydrogen capping



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### A multiple timestep scheme

Relevant dynamical changes in the QM and MM systems typically occur on quite disparate time scales. Accessible time-scale for QM dynamics is the ps time scale, whereas large amplitude conformational changes in large biomolecules occurs on the ns timescale.

- Multiple time step scheme
- oversampling artificially decreasing the MM masses, in this case the generated dynamics is fictious.

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### The CPMD/ GROMOS Interface

Full Hamiltonian QM/MM coupling. The QM/MM interface<sup>3</sup> is tailored to study dynamics of complex biomolecular systems and chemical reactions. The non-bonded part of the total energy is written as:

$$E_{nb}^{QM-MM} = E_{es}^{QM-MM} + E_{steric}^{QM-MM}$$

$$= \sum_{I \in MM} q_I \int \frac{n(\mathbf{r})}{|\mathbf{r} - \mathbf{R}_I|} d\mathbf{r} + \sum_{I \in MM} \sum_{\alpha \in QM} v_{vdW} (|\mathbf{R}_{\alpha} - \mathbf{R}_{\alpha}|)$$
(12)

Short range electrostatic: the spill out problem.

$$E_{es}^{QM-MM} = \sum_{I \in MM} q_I \int n(\mathbf{r}) v_I^{eff}(|\mathbf{r} - \mathbf{R}_I|) d\mathbf{r}$$
(14)

$$v_l^{eff}(r_l) = \frac{r_{cl}^m - r_l^m}{r_{cl}^{m+1} - r_l^{m+1}} \to \frac{1}{r_l}(r_l \to \infty)$$
 (15)

where m-=4 and  $R_{c,I}$  covalent radius for the *I*th atom is a reasonable choice. Laio, VandeVondele, Rothlisberger, J. Chem. Phys. 116, 6941 (2002).

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#### Long Range Electrostatic

The Coulombic electrostatic field is included exactly only for a set of MM atoms in the vicinity of the QM system.

The electrostatic field on the MM atoms that do not belong to the NN set is calculated by a multi-polar expansion

$$\begin{split} \int dr \frac{\rho(r)}{|r-r_j|} &= C \frac{1}{|r_j - \vec{r}|} + \sum_{\alpha} D^{\alpha} \frac{(r_j^{\alpha} - \vec{r}^{\alpha})}{|r_j - \vec{r}|^3} \qquad C = \int dr \rho(r), \\ &+ \frac{1}{2} \sum_{\alpha\beta} Q^{\alpha\beta} \frac{(r_j^{\alpha} - \vec{r}^{\alpha})(r_j^{\beta} - \vec{r}^{\beta})}{|r_j - \vec{r}|^5} \qquad D^{\alpha} = \int dr \rho(r)(r^{\alpha} - \vec{r}^{\alpha}), \\ &+ \mathcal{O} \bigg( \frac{R_q}{|r_j - \vec{r}|} \bigg)^4, \qquad Q^{\alpha\beta} = \int dr \rho(r) [3(r^{\alpha} - \vec{r}^{\alpha})(r^{\beta} - \vec{r}^{\beta}) \\ &- \delta^{\alpha\beta} |r - \vec{r}|^2], \end{split}$$

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### **Bonded Interactions**

Two option have been employed:

- Hydrogen capping
- Carbon monovalent pseudopotentials



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## CPMD/GROMS: validation

QM water in MM (SPC) water



Pair correlation functions for gm/mm compared with continuous line with circles (SPC water) continuous line with diamonds (full QM) Continuous line:  $r_c$  for oxygen and hydrogen are 0.8 and 0.4 Å, respectively; dashed line: r<sub>c</sub> for oxygen and hydrogen are 0.8 and 0.8 Å; dotted line:  $r_c$  for oxygen and hydrogen are 1.4 and 0.8 Å. 1) QM oxygen-MM hydrogen. 2) QM hydrogen-MM oxygen. 3) QM oxygen-MM oxygen.

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#### CPMD/GROMS: validation



Fictitious electronic kinetic energy for a QM water in a box of classical SPC water at 300 K and normal density. The time step is 6 a.u. Continuous line: no modification of the Coulomb potential . Dotted line: modified Coulomb

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Hydrogen-water oxygen pair correlation function for a ammonium ion in a box of water. Continuous line: full quantum result Dashed line: QM/MM result. The quantum ammonium ion is solvated in a box of SPC water.

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# Real space multi-grid approach

This approach<sup>4</sup> is implemented in CP2K<sup>5</sup>. Which is a GPWs code, suing a mixed Gaussian and plane wave basis set for the expansion of the wavefunction and of the charge density.

 This approach is based on the use of a multi-grid technique in conjunction with a Gaussian expansion of the electrostatic potential (GEEP)

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<sup>&</sup>lt;sup>4</sup>JCTC 2005, 1, 1176.

<sup>&</sup>lt;sup>5</sup>www.cp2k.berlios.de

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The exact Coulomb potential is replaced by

$$v_{\mathrm{a}}(\mathbf{r},\mathbf{r}_{\mathrm{a}}) = \frac{\mathrm{Erf}\left(\frac{|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|}{r_{\mathrm{c},\mathrm{a}}}\right)}{|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|}$$

which is the exact potential energy function originated by a Gaussian charge distribution and has the desired property of tending to 1/r at large distances and going smoothly to a constant for small r. This potential can be written as sum of Gaussian functions with different cutoffs according to

$$v_{\mathrm{a}}(\mathbf{r},\mathbf{r}_{\mathrm{a}}) = \frac{\mathrm{Erf}\left(\frac{|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|}{r_{\mathrm{c,a}}}\right)}{|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|} = \sum_{N_{\mathrm{g}}} A_{\mathrm{g}} \exp\left(-\left(\frac{|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|}{G_{\mathrm{g}}}\right)^{2}\right) + R_{\mathrm{low}}(|\mathbf{r}-\mathbf{r}_{\mathrm{a}}|)$$

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The advantage of this decomposition scheme is that grids of different spacing can be used to represent the different contributions to  $v_a(r, r_a)$ .



Sharp Gaussians require fine grids, while coarser grids are necessary for the smoothest components. Gaussians can be truncated beyond a certain threshold value, which makes the collocation of the Gaussians on the grid very efficient.

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Gaussian will be represented on the same number of grid points irrespective of its width. (e.g. a mesh of  $25 \times 25 \times 25$  suffices for an optimal Gaussian representation). The Gaussian can be considered a compact domain function, i.e., it is zero beyond a certain distance. Thus only MM atoms embedded into the QM box, or close to it, will contribute to the finest grid levels.

The speed up is abut 2 order of magnitude

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CP-PAW/COSMO Interface

#### Embedding in continuum environment

 $\ensuremath{\mathsf{CP}}\xspace{\mathsf{PAW}}\xspace/\mathsf{COSMO}\xspace:$  the environment is treated with a conductor-like screening model

- homogeneous and inert solvent treated as dielectric continuum: suitable if environment is not involved in the ongoing chemistry
- the essential aspect of this approach is to treat the surface charges at the cavity boundaries as fictious dynamical variables:

$$\mathcal{L}^{QM/MM} = \mathcal{L}^{QM}_{CP} + \mathcal{L}^{MM} + \mathcal{L}^{COSMO}$$
(16)  
$$\mathcal{L}^{COSMO} = \sum_{I} \frac{1}{2} M_{I} \dot{Q}_{I}^{2} + G^{COSMO}_{es} (\{Q_{I}\}, \{\mathbf{R}_{\alpha}\})$$
$$- G_{steric} (\{\mathbf{R}_{\alpha}\}) - \sum_{I} k(1 - \Theta_{I}) Q_{I}^{2}$$
(17)

where the

$$G_{steric} = \sigma_0 + \sigma_1 A \tag{18}$$

with A surface area of the cavity and  $\sigma_i$  empirical fixed parameters. Marialore Sulpizi Density Functional Theory: from theory to Applications

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#### CP-PAW/COSMO Interface



COSMO surface of a pentaacrylate molecule (red = negative, green = positive equilibrium layer).

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- A bonded term between the QM and MM parts,  $E_b^{QM-MM}$  does not exists in this approach.
- The fictious variables Q<sub>I</sub>, with associated inertia parameters M<sub>I</sub><sup>Q</sup>, are discretized, scaled surface charges, located at the center of segments s<sub>I</sub>.
- Switching functions are introduced to switch off the charges of the segments which are not exposed to the solvent.
- Last term in eq:17 is a penalty function introduced to keep the instantaneous switched off charges from blowing-up during the dynamics

CP-PAW/COSMO Interface

The COSMO free energy can be written down in terms of of electrostatic solute-solvent, solvent-solvent, and self energy contributions:

$$G_{es}^{COSMO} = \sum_{I} Q_{I} \Theta_{I} \int_{V} \frac{n(\mathbf{r})}{|\mathbf{r} - \mathbf{s}_{I}|} d\mathbf{r} + \frac{1}{f} \sum_{I < J} \frac{Q_{I} \Theta_{I} Q_{J} \Theta_{J}}{|\mathbf{s}_{I} - \mathbf{s}_{J}|} \quad (19)$$
$$+ \frac{c}{f} \sum_{I} \frac{Q_{I}^{2} \Theta_{I}^{2}}{\sqrt{a_{I}}} \quad (20)$$
$$G_{steric} = \sigma_{0} + \sigma_{1} \sum_{I} a_{I} \Theta_{I} \quad (21)$$

f is a constant screening factor, c is a geometry-dependent constant and V is the cavity volume. Spill-out is prevented using the Gaussian model density as described for the CP-PAW/AMBER.

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### A QM/MM application: reaction mechanism of caspase

Caspase-3 is a cysteine protease involved in the programmed cell-death (apoptosis). Idea: study the enzymatic reaction in the protein environment. Calculate free energy barrier associated with the proposed mechanism<sup>6</sup>.



<sup>6</sup>Sulpizi *et al*, Proteins, 52(2): 212 (2003). Marialore Sulpizi Density Functional Theory: from theory to Applications





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The free energy barrier associated with this step is only 5 kcal/mol, so the previous step results to be the rate determing one.

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#### Importance of the environment

- No secondary structure elements which are relevant to active site electric field (as i.e. in papain and cathepsins)
- Protein structure seems to act as geometrical constraint which reduces entropy of reaction provides proper conformation to catalytic dyad

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