

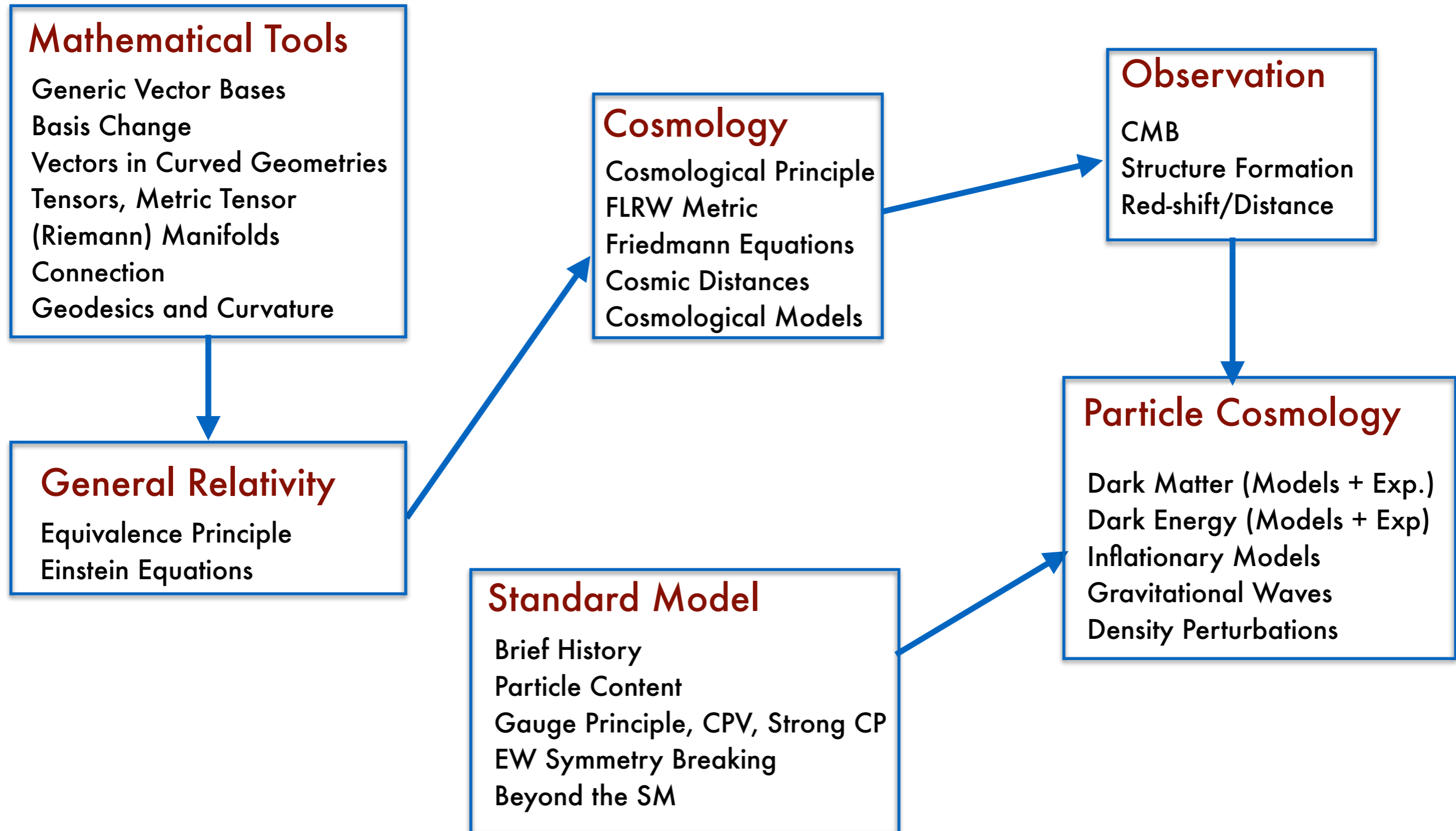
Introductory Particle Cosmology

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Lecture 4



Datum	Von	Bis	Raum
1 Di, 17. Apr. 2018	10:00	12:00	05 119 Minkowski-Raum
2 Do, 19. Apr. 2018	08:00	10:00	05 119 Minkowski-Raum
3 Di, 24. Apr. 2018	10:00	12:00	05 119 Minkowski-Raum
4 Do, 26. Apr. 2018	08:00	10:00	05 119 Minkowski-Raum
5 Do, 3. Mai 2018	08:00	10:00	05 119 Minkowski-Raum
6 Di, 8. Mai 2018	10:00	12:00	05 119 Minkowski-Raum
7 Di, 15. Mai 2018	10:00	12:00	05 119 Minkowski-Raum
8 Do, 17. Mai 2018	08:00	10:00	05 119 Minkowski-Raum
9 Di, 22. Mai 2018	10:00	12:00	05 119 Minkowski-Raum
10 Do, 24. Mai 2018	08:00	10:00	05 119 Minkowski-Raum
11 Di, 29. Mai 2018	10:00	12:00	05 119 Minkowski-Raum
12 Di, 5. Jun. 2018	10:00	12:00	05 119 Minkowski-Raum
13 Do, 7. Jun. 2018	08:00	10:00	05 119 Minkowski-Raum
14 Di, 12. Jun. 2018	10:00	12:00	05 119 Minkowski-Raum
15 Do, 14. Jun. 2018	08:00	10:00	05 119 Minkowski-Raum
16 Di, 19. Jun. 2018	10:00	12:00	05 119 Minkowski-Raum
17 Do, 21. Jun. 2018	08:00	10:00	05 119 Minkowski-Raum
18 Di, 26. Jun. 2018	10:00	12:00	05 119 Minkowski-Raum
19 Do, 28. Jun. 2018	08:00	10:00	05 119 Minkowski-Raum
20 Di, 3. Jul. 2018	10:00	12:00	05 119 Minkowski-Raum
21 Do, 5. Jul. 2018	08:00	10:00	05 119 Minkowski-Raum



H. Minkowski
(1864-1909)

Physics Dept. Building, 5th Floor

The Cosmological Principle

The FLRW metric

Friedmann Equations

Cosmological Models and Parameters

LambdaCDM

The Universe is spatially homogeneous and isotropic

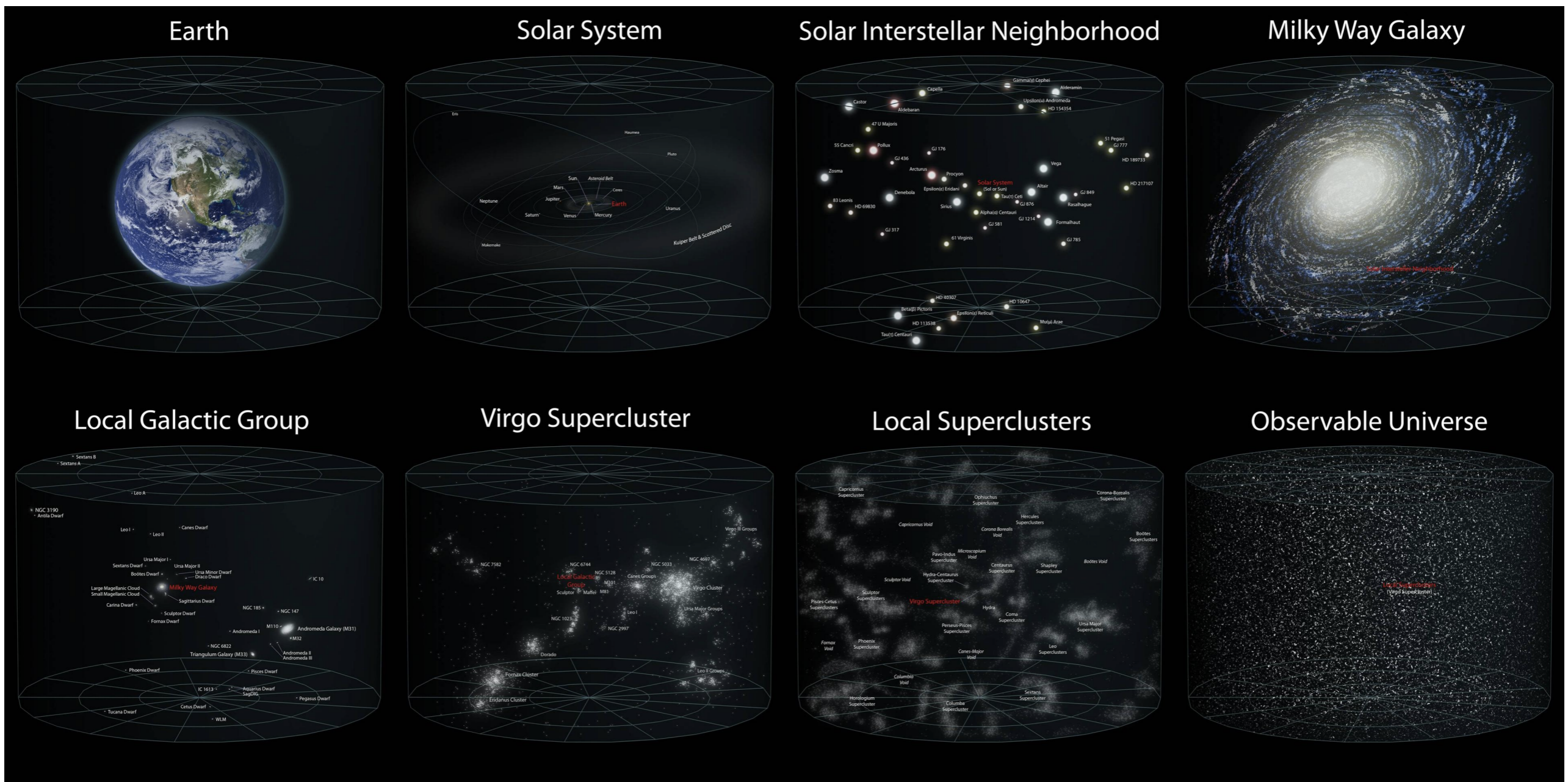


Image from Cryhavoc

From the Cosmological Principle

we obtain the **Friedmann-Lemaitre-Robertson-Walker** metric:

$$ds^2 = dt^2 - R^2(t) \left[\frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right]$$

Noting the rescaling invariance

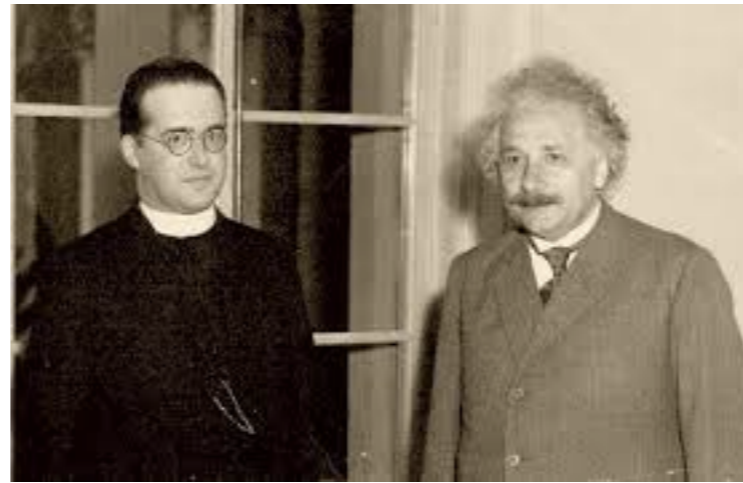
$$\begin{aligned} R &\rightarrow \frac{R}{\lambda} \\ r &\rightarrow \lambda r \\ k &\rightarrow \frac{k}{\lambda^2} \end{aligned}$$

we can always choose $k=+1,0,-1$.

R can be scaled to e.g. 1 for $t=\text{today}$. Usually $R(t)$ is called $a(t)$: the scale factor.



Alexander Friedmann (1888-1925)
Russian Physicist and Mathematician



Georges Lemaître (1894-1966)
Belgian Catholic Priest and Astronomer

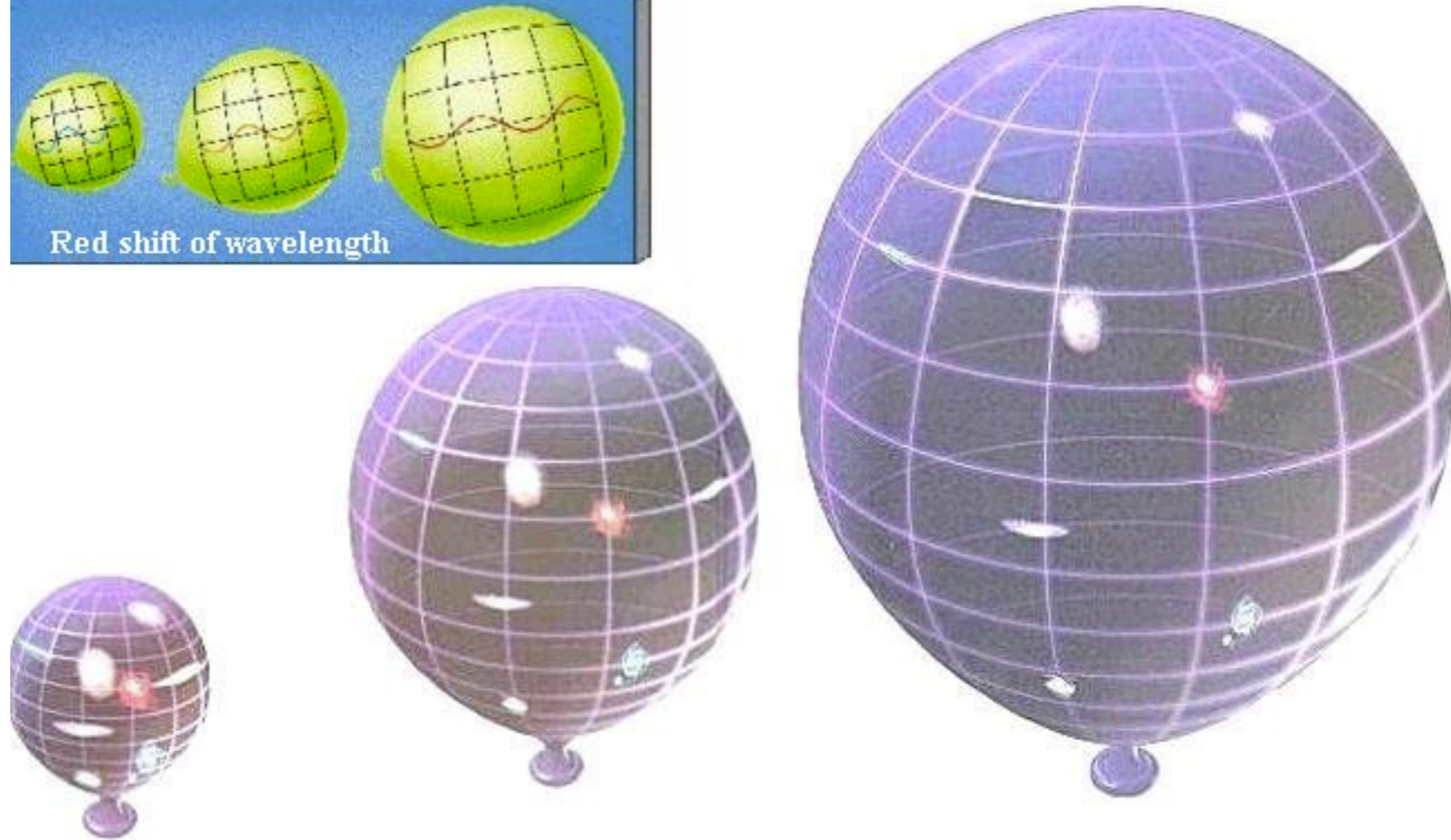
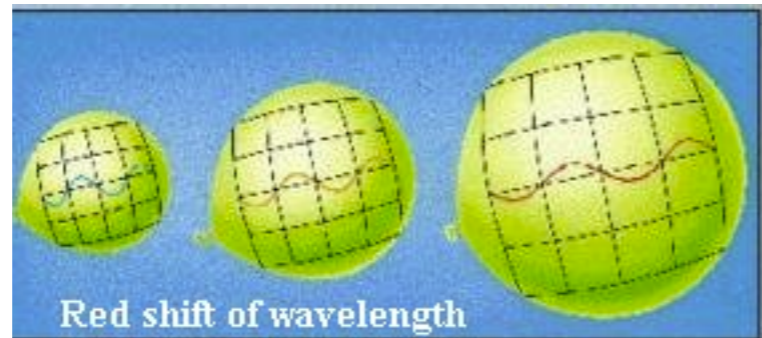


Howard P. Robertson (1903-1961)
USA physicist and mathematician



Arthur G. Walker (1909-2001)
UK Mathematician

Comoving Coordinates



Non-zero components of the FLRW metric tensor

$$\begin{aligned}
 g_{00} &= 1 \\
 g_{11} &= -\frac{a^2(t)r^2}{1-kr^2} \\
 g_{22} &= -a^2(t)r^2 \\
 g_{33} &= -a^2(t)r^2 \sin^2 \theta
 \end{aligned}$$

Non-zero components of the FLRW Christoffel symbols

$$\begin{aligned}
 \Gamma_{11}^0 &= \frac{\dot{a}a}{1-kr^2} \quad ; \quad \Gamma_{22}^0 = a\dot{a}r^2 \quad ; \quad \Gamma_{33}^0 = a\dot{a}r^2 \sin^2 \theta \\
 \Gamma_{01}^1 &= \Gamma_{10}^1 = \Gamma_{02}^2 = \Gamma_{20}^2 = \Gamma_{03}^3 = \Gamma_{30}^3 = \frac{\dot{a}}{a} \\
 \Gamma_{22}^1 &= -r(1-kr^2) \quad ; \quad \Gamma_{33}^1 = -r(1-kr^2) \sin^2 \theta \\
 \Gamma_{12}^2 &= \Gamma_{21}^2 = \Gamma_{13}^3 = \Gamma_{31}^3 = \frac{1}{r} \\
 \Gamma_{33}^2 &= -\sin \theta \cos \theta \quad ; \quad \Gamma_{23}^3 = \Gamma_{32}^3 = \cot \theta \quad ,
 \end{aligned}$$

Non-zero components of the Ricci tensor

$$\begin{aligned}
 R_{00} &= -3\frac{\ddot{a}}{a} \\
 R_{11} &= \frac{a\ddot{a} + 2\dot{a}^2 + 2k}{1-kr^2} \\
 R_{22} &= r^2(a\ddot{a} + 2\dot{a}^2 + 2kr) \\
 R_{33} &= r^2(a\ddot{a} + 2\dot{a}^2 + 2kr) \sin^2 \theta
 \end{aligned}$$

Ricci curvature scalar

$$R = \frac{6}{a^2}(a\ddot{a} + 2\dot{a}^2 + k)$$

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda g_{\mu\nu} = 8\pi T_{\mu\nu}$$

Calculated in the previous slide

Still missing...but we can use again the cosmological principle:

E/m tensor for an isotropic homogeneous "fluid":

$$T_{\mu\nu} = (\rho + P)v_{\mu}v_{\nu} - Pg_{\mu\nu}$$

Density Pressure 4-velocity

The "fluid" is at rest wrt the co-moving coordinates:
 $v=(1,0,0,0)$

Note this!

GR is in general "hard" to solve:
 the metric shows up on BOTH sides of the Einstein Equations..

Putting together the previous calculations, we can obtain the **Friedmann equations**:

Acceleration equation:

Hubble par. equation:

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P)$$
$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2}$$

The Friedmann equations are the Einstein equations for the FLRW metric and an isotropic homogenous fluid.

How to solve them for the scale factor $a(t)$? We need to know pressure and density of the “fluid”, or at least a relation between them: an **equation of state**.

Energy-momentum conservation: $\nabla_{\mu} T_{\nu}^{\mu} = \partial_{\mu} T_{\nu}^{\mu} + \Gamma_{\mu\beta}^{\mu} T_{\nu}^{\beta} - \Gamma_{\mu\nu}^{\beta} T_{\beta}^{\mu} = 0$

For the 0-th component: $\partial_0 \rho(t) + 3 \frac{\dot{a}(t)}{a(t)} (\rho(t) + P(t)) = 0$

Choose the generic equation of state: $P = w\rho$

$$\frac{\dot{\rho}}{\rho} = -3(1+w) \frac{\dot{a}}{a}$$

$$\rho(t) \propto a(t)^{-3(1+w)}$$

$$H(t) = \frac{1}{a} \frac{da}{dt} \quad \text{Hubble Parameter}$$

$$q(t) = -\frac{1}{a} \frac{d^2a}{dt^2} / \left(\frac{1}{a} \frac{da}{dt} \right)^2 \quad \text{Deceleration}$$

$$j(t) = \frac{1}{a} \frac{d^3a}{dt^3} / \left(\frac{1}{a} \frac{da}{dt} \right)^3 \quad \text{"Jerk"}$$

$$s(t) = \frac{1}{a} \frac{d^4a}{dt^4} / \left(\frac{1}{a} \frac{da}{dt} \right)^4 \quad \text{"Snap" (or Jounce)}$$

$$c(t) = \frac{1}{a} \frac{d^5a}{dt^5} / \left(\frac{1}{a} \frac{da}{dt} \right)^5 \quad \text{"Crackle"}$$

$$p(t) = \frac{1}{a} \frac{d^6a}{dt^6} / \left(\frac{1}{a} \frac{da}{dt} \right)^6 \quad \text{"Pop"}$$

...

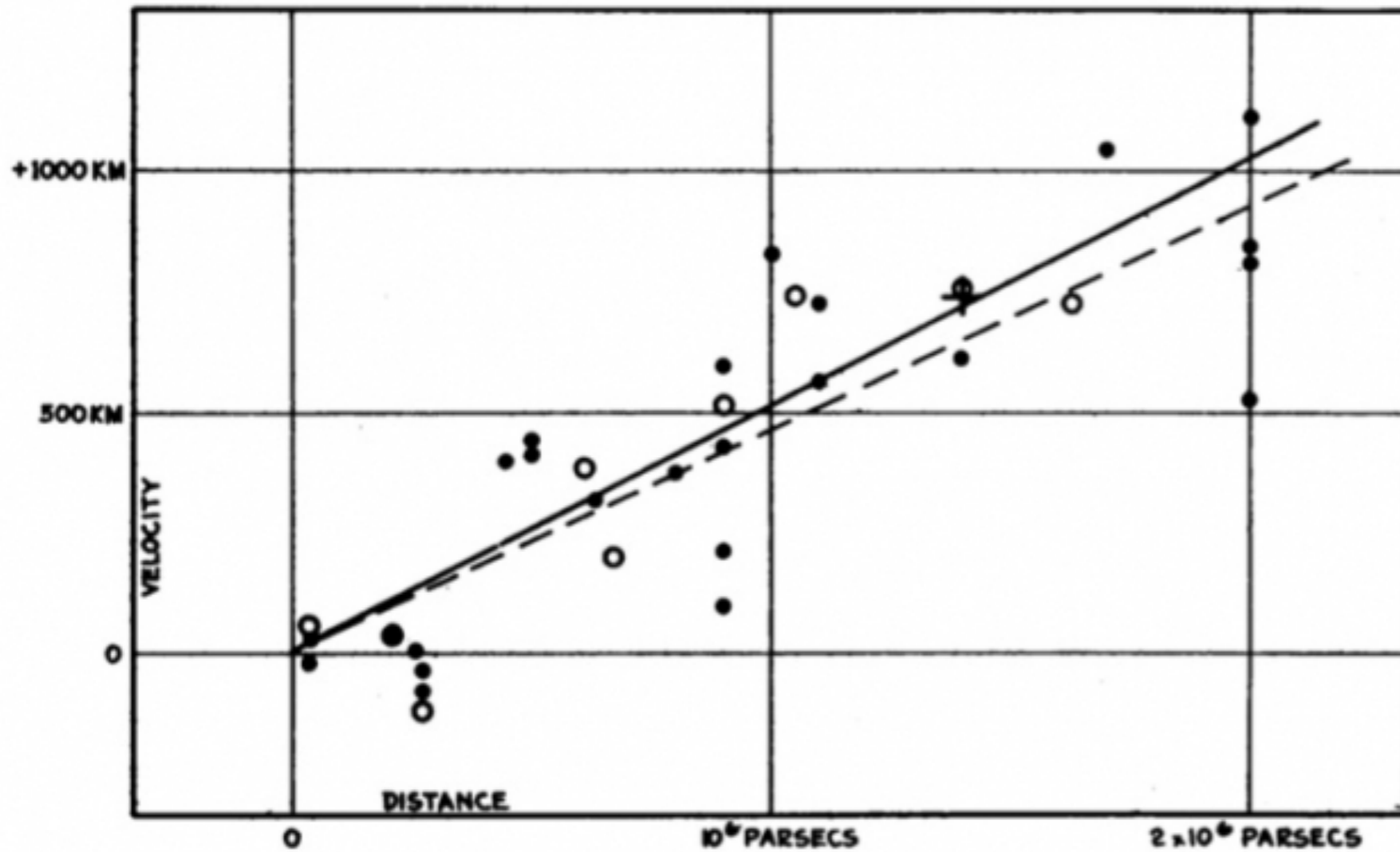
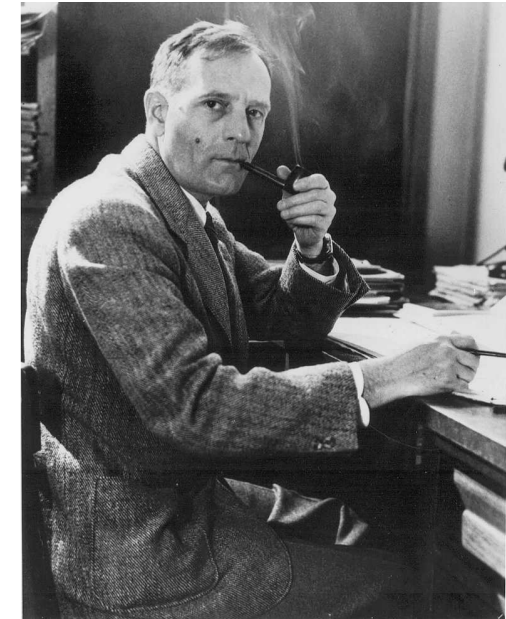


FIGURE 1

Velocity-Distance Relation among Extra-Galactic Nebulae.



Edwin Hubble (1889 - 1953)

$$d \propto v \quad \Rightarrow \quad \frac{\dot{a}}{a} = H \quad \Rightarrow \quad \dot{a} = H a$$

Friedmann Equations
with Cosmological Constant Term

$$\frac{\ddot{a}}{a} = \dot{H} + H^2 = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3}$$

$$H^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3}.$$

Equations of state
 $P = w\rho$

$$w = 0 \Rightarrow \rho \sim \frac{1}{a^3} \quad \text{"Dust"}$$

$$w = 1/3 \Rightarrow \rho \sim \frac{1}{a^4} \quad \text{"Radiation"}$$

$$w = -1 \Rightarrow \rho \sim \text{const.} \quad \text{"Vacuum"}$$

“Dust”: $w = 0 \Rightarrow \rho \cdot a(t)^3 = A$

Define the **Conformal Time** $\frac{d\eta}{dt} = \frac{1}{a}$

The 2nd Friedmann equation becomes $a'^2 = \frac{8\pi G}{3} Aa - ka^2$

With solutions

$$\begin{aligned}
 k = 1 &\Rightarrow a = \frac{4\pi GA}{3}(1 - \cos \eta) & ; & t = \frac{4\pi GA}{3}(\eta - \sin \eta) \\
 k = 0 &\Rightarrow a = \frac{2\pi GA}{3}\eta^2 & ; & t = \frac{2\pi GA}{9}\eta^3 \\
 k = -1 &\Rightarrow a = \frac{4\pi GA}{3}(\cosh \eta - 1) & ; & t = \frac{4\pi GA}{3}(\sinh \eta - \eta)
 \end{aligned}$$

Radiation: $w = 1/3 \Rightarrow \dot{a}^2 = \frac{8\pi G A}{3 a^2} - k$

With solutions $k = 1 \Rightarrow a = \sqrt{2\sqrt{\frac{8\pi G A}{3 a^2}} t - t^2}$

$$k = 0 \Rightarrow a = \sqrt{2\sqrt{\frac{8\pi G A}{3 a^2}} t}$$

$$k = -1 \Rightarrow a = \sqrt{2\sqrt{\frac{8\pi G A}{3 a^2}} t + t^2}$$

Vacuum: $P = \rho = 0 \Rightarrow \dot{a}^2 = \frac{\Lambda a^2}{3} - k$

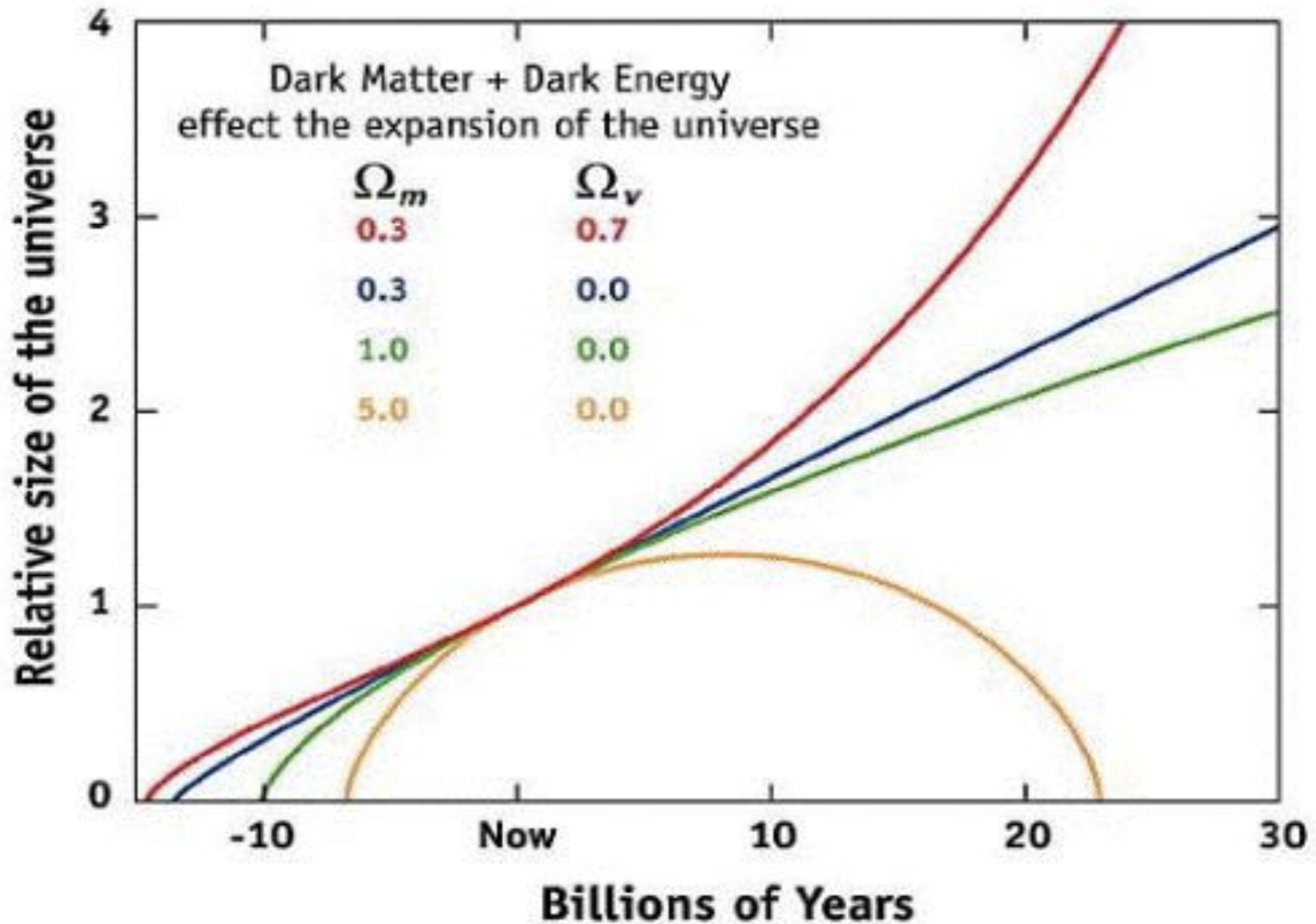
With solutions

$$k = 1 \Rightarrow a = \sqrt{\frac{3}{\Lambda}} \cosh \left(\sqrt{\frac{\Lambda}{3}} t \right)$$

$$k = 0 \Rightarrow a = \sqrt{\frac{3}{\Lambda}} \exp \left(\sqrt{\frac{\Lambda}{3}} t \right)$$

← Note this result

$$k = -1 \Rightarrow a = \sqrt{\frac{3}{\Lambda}} \sinh \left(\sqrt{\frac{\Lambda}{3}} t \right)$$



Remember the Hubble parameter: $H(t) = \frac{\dot{a}(t)}{a(t)}$

The 2nd Friedmann Equation becomes $H^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} \Rightarrow \frac{8\pi G}{3H^2}\rho - 1 = \frac{k}{H^2a^2}$

Introducing the **Critical Density** $\rho_c = \frac{3H^2}{8\pi G}$

And the **Density Parameter** $\Omega = \rho/\rho_c$ we have: $\Omega - 1 = \frac{k}{H^2a^2}$

$$\rho < \rho_c \iff \Omega < 1 \iff k = -1 \quad (\text{Open})$$

and therefore $\rho = \rho_c \iff \Omega = 1 \iff k = 0 \quad (\text{Flat})$

$$\rho > \rho_c \iff \Omega > 1 \iff k = 1 \quad (\text{Closed})$$

If more than a component is present $\rho_{TOT}(a) = \sum_i \rho_i(a) = \rho_C \sum_i \Omega_i a^{-3(1+w_1)}$

The Friedmann equation is then $\frac{k}{a^2} = H^2(\Omega_{TOT} - 1)$

And introducing the today's observed density parameters

$$\frac{k}{a_0^2} = H_0^2(\Omega_m + \Omega_r + \Omega_\Lambda - 1)$$

