

# Quantum Field Theory

## Exercise 11

January 26, 2017

-to be handed in by 02.02.2017 (12:00 h) to the letterbox No. 37 (“relativistische QFT”) in the foyer of Staudingerweg 7.

**(1) Path integrals in Quantum Mechanics (45 points)**

In the lecture you derived the amplitude for the propagation of a particle from  $q$  at time 0 to  $q'$  at time  $t$  to be

$$\langle q'(t) | e^{-iHt} | q(0) \rangle = \int \prod_{k=1}^n dq_k \prod_{j=0}^n \frac{dp_j}{2\pi} e^{ip_j(q_{j+1}-q_j)} e^{-i\frac{p_j^2}{2m}\delta t}, \quad (1)$$

(2)

where the potential term is now set to zero for simplicity. Useful formulae are

$$\dot{q}_j = \frac{q_{j+1} - q_j}{\delta t}, \quad \delta t = \frac{t}{n+1}. \quad (3)$$

- (a) **(20 points)** Consider the general  $p_j$  integral. Complete the square in order to make it gaussian. When solving the integral treat it as a real integral, in particular use  $\int_{-\infty}^{\infty} e^{-cx^2} dx = (\frac{\pi}{c})^{\frac{1}{2}}$  where  $c$  can contain imaginary unit. Using the result of the integration prove

$$\int \prod_{j=0}^n \frac{dp_j}{2\pi} e^{ip_j(q_{j+1}-q_j)} e^{-i\frac{p_j^2}{2m}\delta t} = \left(\frac{m}{2\pi i \delta t}\right)^{\frac{n+1}{2}} \exp\left[\frac{im}{2\delta t} \sum_{j=0}^n (q_{j+1} - q_j)^2\right]. \quad (4)$$

- (b) **(25 points)** Now the total amplitude equals

$$\left(\frac{m}{2\pi i \delta t}\right)^{\frac{n+1}{2}} \int \prod_{k=1}^n dq_k \exp\left[\frac{im}{2\delta t} \sum_{j=0}^n (q_{j+1} - q_j)^2\right]. \quad (5)$$

Try to solve also the integrals over  $dq_k$ . First integrate over  $q_1$ , then  $q_2$ , etc. and try to look for a pattern. Again, as in part (1 a), complete the square, treat the integrals as real and apply the formula for gaussian integrals.

*Hint : The result is*

$$\int \prod_{k=1}^n dq_k \exp\left[\frac{im}{2\delta t} \sum_{j=0}^n (q_{j+1} - q_j)^2\right] = \left[\frac{2i\pi\delta t}{m}\right]^{\frac{n}{2}} \sqrt{\frac{n!}{(n+1)!}} \exp\left[\frac{im(q_{n+1} - q_0)^2}{2(n+1)\delta t}\right].$$

Since  $q_0 = q$  and  $q_{n+1} = q'$  it is now easy to express this result, together with a prefactor from (5) (and using eqs. (3)) as

$$\langle q'(t) | e^{-iHt} | q(0) \rangle = \sqrt{\frac{m}{2\pi it}} \exp \left[ \frac{im(q' - q)^2}{2t} \right],$$

which depends only on the initial and final position, the time and the mass of the particle.

**(2) Path Integral for a Free Complex Scalar Field (35 points)**

Obtain the following form of a path integral  $Z_0(J, J^*)$  for a free complex scalar field

$$Z_0(J, J^*) = \text{Exp} \left[ - \int d^4x \int d^4x' J^*(x') \Delta(x - x') J(x) \right], \quad (6)$$

where  $\Delta(x - x')$  is the Feynman propagator in position space.

*Hint: Note that in the case of a complex scalar field one has to introduce two source terms of the form  $J^*(x)\phi(x) + J(x)\phi^*(x)$ .*

**(3) Polarized  $e^+e^- \rightarrow \mu^+\mu^-$  Cross Section (20 points)**

Starting from the matrix element for  $e^+e^- \rightarrow \mu^+\mu^-$  process

$$i\mathcal{M}(e^-(p)e^+(p') \rightarrow \mu^-(k)\mu^+(k')) = \frac{ie^2}{(p + p')^2} (\bar{v}(p')\gamma^\mu u(p)) (\bar{u}(k)\gamma_\mu v(k')) \quad (7)$$

calculate the differential cross section

$$\frac{d\sigma}{d\Omega}(e_R^+e_L^- \rightarrow \mu_L^+\mu_R^-), \quad (8)$$

where the indices L and R label left and right handed particles, respectively. Work in the massless limit ( $m_e, m_\mu \rightarrow 0$ ). Express the differential cross section in terms of the incoming electron energy  $E$  and the angle between the incoming electron and the outgoing muon.