

Quantum Field Theory

Exercise 3

November 10, 2016

-to be handed in by 17.11.2016 (12:00 h) to the letterbox No. 37 (“relativistische QFT”) in the foyer of Staudingerweg 7.

1. Relativistic particle in a homogeneous magnetic field (35 points)

The coupling of the Dirac equation to an electromagnetic field is accomplished by the substitution $i\partial_\mu \rightarrow i\partial_\mu - eA_\mu$, where $A^\mu = (\Phi, \mathbf{A})$ is the electromagnetic 4-potential. The Dirac equation in presence of an electromagnetic field then reads:

$$((i\cancel{\partial} - e\cancel{A}) - m)\psi = 0 \quad (1)$$

(a) (5 points) Show that eq. 1 can be written as

$$\hat{E}\psi = \boldsymbol{\alpha}(\hat{\mathbf{p}} - e\mathbf{A})\psi + e\Phi\psi + m\gamma^0\psi, \quad (2)$$

with $\alpha^i = \gamma^0\gamma^i = \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}$ and the usual quantum-mechanical replacements $\hat{p}^0 = \hat{p}_0 = \hat{E} = i\partial_t$ and $\hat{\mathbf{p}} = -i\nabla$. Note that here we use the *Dirac basis*, where $\gamma_0 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \beta$.

Consider now a particle in a homogeneous magnetic field \mathbf{B} . Use the gauge $A^0 = A^2 = A^3 = 0$, $A^1 = By$.

(b) (5 points) Show that eq. 2 now reads

$$(\hat{E} - m)\phi = \boldsymbol{\sigma}(\hat{\mathbf{p}} - eBy\hat{\mathbf{e}}_1)\chi \quad (3)$$

$$(\hat{E} + m)\chi = \boldsymbol{\sigma}(\hat{\mathbf{p}} - eBy\hat{\mathbf{e}}_1)\phi \quad (4)$$

for the two components ϕ, χ of $\psi = \begin{pmatrix} \phi \\ \chi \end{pmatrix}$.

(c) (15 points) Derive the following equation for ϕ :

$$(\hat{E}^2 - m^2)\phi = (\hat{p}_y^2 + \hat{p}_z^2 + (\hat{p}_x - eBy)^2 + eB\sigma^3)\phi \quad (5)$$

$$= (\hat{p}_y^2 + \hat{p}_z^2 + (\hat{p}_x - eBy)^2 + 2eB\hat{s}_z)\phi \quad (6)$$

Since the right hand side of eq. 5 commutes with \hat{p}_x, \hat{p}_z and \hat{s}_z , one can use the ansatz $\phi = e^{i(p_x x + p_z z)}\omega(p_y, y)\rho(s_z)$ for the stationary solution.

(d) **(10 points)** Derive a differential equation for $\omega(p_y, y)$. What does it remind you of? Deduce the corresponding energy levels for a relativistic particle in a homogeneous magnetic field.

2. Gordon Identity (25 points)

Derive the so-called *Gordon identity*

$$\bar{u}(p')\gamma^\mu u(p) = \bar{u}(p') \left[\frac{P^\mu}{2M} + \frac{i\sigma^{\mu\nu}q_\nu}{2M} \right] u(p),$$

where $P = p' + p$, $q = p' - p$ and $\sigma^{\mu\nu} = \frac{i}{2} [\gamma^\mu, \gamma^\nu]$.

3. Projection Operators (20 points)

Show that $P_\pm = \pm \frac{\not{p} \pm m}{2m}$ represent a complete sets of projection operators, *i.e.* satisfy the conditions

$$P_i P_j = \delta_{i,j} P_i, \quad \sum_i P_i = 1.$$

In addition, show that P_\pm are the projection operators on positive and negative energy solutions (u and v spinors) for arbitrary particle momentum.

4. γ Matrices (maximum 20 points): you are welcome to do as many as you like, but only the first to give 20 points will be marked.

Without using an explicit representation for the γ matrices show that:

- (a) **(2.5 points)** $\gamma_\mu \gamma^\mu = 4$
- (b) **(2.5 points)** $\text{Tr}[\gamma^\mu \gamma^\nu] = 4g^{\mu\nu}$
- (c) **(5 points)** $\text{Tr}[\not{a} \not{b} \not{c} \not{d}] = 4[(a \cdot b)(c \cdot d) - (a \cdot c)(b \cdot d) + (a \cdot d)(b \cdot c)]$
- (d) **(5 points)** $\gamma_5 = \frac{i}{4!} \varepsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma$
- (e) **(5 points)** $\gamma_5^2 = \mathbb{1}$
- (f) **(5 points)** $\{\gamma_5, \gamma^\mu\} = 0$
- (g) **(5 points)** $\text{Tr}[\gamma_5] = 0$
- (h) **(5 points)** $\text{Tr}[\gamma^\mu \gamma^\nu \gamma_5] = 0$
- (i) **(5 points)** $\text{Tr}[\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma_5] = 4i \varepsilon^{\mu\nu\rho\sigma}$
- (j) **(5 points)** $\text{Tr}[\gamma^{\mu_1} \dots \gamma^{\mu_n}] = 0$ if n is odd

where $\not{a} \equiv \gamma^\mu a_\mu$, $\gamma_5 \equiv i\gamma^0\gamma^1\gamma^2\gamma^3$ and $\varepsilon_{0123} = +1$