Homogenization of Hyperbolic-type Equations with Periodic Coefficients

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Let Γ be a lattice in \mathbb{R}^d , let Ω be the cell of Γ . Let $\widetilde{\Gamma}$ be the dual lattice. By $\widetilde{\Omega}$ we denote the central Brillouin zone of $\widetilde{\Gamma}$. Example:

$$\Gamma = \mathbb{Z}^d, \ \ \Omega = (0,1)^d; \ \ \ \widetilde{\Gamma} = (2\pi\mathbb{Z})^d, \ \ \widetilde{\Omega} = (-\pi,\pi)^d.$$

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Let $\varepsilon > 0$ be a parameter. We use the notation $\phi^{\varepsilon}(\mathbf{x}) = \phi\left(\frac{\mathbf{x}}{\varepsilon}\right)$.

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Main object

In $L_2(\mathbb{R}^d; \mathbb{C}^n)$, we consider elliptic second order DO

 $A_{\varepsilon} = b(\mathbf{D})^* g^{\varepsilon}(\mathbf{x}) b(\mathbf{D}).$

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Homogenization of hyperbolic equations

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Here $g(\mathbf{x})$ is Hermitian $(m \times m)$ -matrix with complex entries. We assume that $g(\mathbf{x})$ is Γ -periodic, bounded and positive definite:

 $c'\mathbf{1}_m \leqslant g(\mathbf{x}) \leqslant c''\mathbf{1}_m, \quad 0 < c'' \leqslant c'' < \infty.$

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 $b(\mathbf{D}) = \sum_{j=1}^{d} b_j D_j$ is a first order $(m \times n)$ -matrix DO; b_j are constant matrices, and $m \ge n$. The symbol $b(\boldsymbol{\xi}) = \sum_{j=1}^{d} b_j \xi_j$ is such that

 $\operatorname{rank} b(\boldsymbol{\xi}) = n, \quad 0 \neq \boldsymbol{\xi} \in \mathbb{R}^d.$

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Precise Definition: A_{ε} is a selfadjoint operator in $L_2(\mathbb{R}^d; \mathbb{C}^n)$ generated by the quadratic form

$$a_{\varepsilon}[\mathbf{u},\mathbf{u}] = \int_{\mathbb{R}^d} \langle g^{\varepsilon}(\mathbf{x}) b(\mathsf{D}) \mathbf{u}, b(\mathsf{D}) \mathbf{u}
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Under our assumptions,

$$c_0\int_{\mathbb{R}^d}|\mathsf{D}\mathsf{u}|^2d\mathsf{x}\leqslant a_arepsilon[\mathsf{u},\mathsf{u}]\leqslant c_1\int_{\mathbb{R}^d}|\mathsf{D}\mathsf{u}|^2\,d\mathsf{x},\quad\mathsf{u}\in H^1(\mathbb{R}^d;\mathbb{C}^n).$$

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Example: $A_{\varepsilon} = -\operatorname{div} g^{\varepsilon}(\mathbf{x}) \nabla = \mathbf{D}^* g^{\varepsilon}(\mathbf{x}) \mathbf{D}.$

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Problem

The problem is to study the behavior of the operator

 $\cos(\tau A_{\varepsilon}^{1/2}), \ \tau \in \mathbb{R},$

for small ε , and to apply the results to the Cauchy problem for the hyperbolic-type equation:

 $\begin{cases} \partial_{\tau}^{2} \mathbf{u}_{\varepsilon}(\mathbf{x},\tau) = -(A_{\varepsilon}\mathbf{u}_{\varepsilon})(\mathbf{x},\tau), & \mathbf{x} \in \mathbb{R}^{d}, \quad \tau \in \mathbb{R}, \\ \mathbf{u}_{\varepsilon}(\mathbf{x},0) = \boldsymbol{\phi}(\mathbf{x}), & \partial_{\tau}\mathbf{u}_{\varepsilon}(\mathbf{x},0) = \boldsymbol{\psi}(\mathbf{x}). \end{cases}$

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We show that, in some sense,

$$\cos(au A_arepsilon^{1/2})\sim \cos(au(A^0)^{1/2}), \quad arepsilon o 0,$$

where A^0 is the effective operator with constant effective coefficients.

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Definition of the effective matrix:

Let $\Lambda(x)$ be the $(n \times m)$ -matrix-valued Γ -periodic solution of the problem

$$b(\mathbf{D})^*g(\mathbf{x})(b(\mathbf{D})\Lambda(\mathbf{x})+\mathbf{1}_m)=0; \quad \int_\Omega \Lambda(\mathbf{x})\,d\mathbf{x}=0.$$

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Then g^0 is an $(m \times m)$ -matrix given by

$$g^0 = |\Omega|^{-1} \int_{\Omega} \widetilde{g}(\mathsf{x}) d\mathsf{x}, \quad \widetilde{g}(\mathsf{x}) \coloneqq g(\mathsf{x})(b(\mathsf{D})\wedge(\mathsf{x}) + \mathbf{1}_m).$$

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$$\|(A_{\varepsilon}+I)^{-1}-(A^{0}+I)^{-1}\|_{L_{2}(\mathbb{R}^{d})\to L_{2}(\mathbb{R}^{d})}\leqslant C\varepsilon.$$

$$(1)$$

Homogenization of hyperbolic equations

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The method was based on the scaling transformation, the Floquet-Bloch theory and the analytic perturbation theory.

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The method was based on the scaling transformation, the Floquet-Bloch theory and the analytic perturbation theory.

In 2004 T. Suslina proved that

$$\|e^{-A_{\varepsilon}\tau}-e^{-A^{0}\tau}\|_{L_{2}(\mathbb{R}^{d})\to L_{2}(\mathbb{R}^{d})}\leqslant C(\tau)\varepsilon,\quad \tau>0.$$
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A different approach to operator error estimates was suggested by
 V. Zhikov and S. Pastukhova (2005–2006).

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 $\|\exp(-i\tau A_{\varepsilon}) - \exp(-i\tau A^{0})\|_{H^{3}(\mathbb{R}^{d}) \to L_{2}(\mathbb{R}^{d})} \leq (\widetilde{C}_{1} + \widetilde{C}_{2}|\tau|)\varepsilon,$ $\|\cos(\tau A_{\varepsilon}^{1/2}) - \cos(\tau (A^{0})^{1/2})\|_{H^{2}(\mathbb{R}^{d}) \to L_{2}(\mathbb{R}^{d})} \leq (C_{1} + C_{2}|\tau|)\varepsilon.$ (3)

Main questions

Is the result

$$\|\cos(\tau A_{\varepsilon}^{1/2}) - \cos(\tau (A^0)^{1/2})\|_{H^2(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant C(\tau)\varepsilon$$

sharp (with respect to the type of operator norm)?

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Answers

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Is it possible to improve this result, probably, under some additional assumptions?

Answers

YES!, YES!

Similar results were obtained also for nonstationary Schrödinger type equations by **T.A. Suslina**.

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Reduction 1: Scaling transformation

Question: for what (minimal) s the estimate

$$\|e^{-i\tau A_{\varepsilon}^{1/2}} - e^{-i\tau (A^0)^{1/2}}\|_{H^s(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant C(\tau)\varepsilon$$
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holds? Let $H_0 = -\Delta$. Clearly, (4) is equivalent to

$$\left\| \left(e^{-i\tau A_{\varepsilon}^{1/2}} - e^{-i\tau (A^0)^{1/2}} \right) (H_0 + I)^{-s/2} \right\|_{L_2(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant C(\tau) \varepsilon.$$
 (5)

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(5)

Next, by the scaling transformation, (5) is equivalent to

$$\left\| \left(e^{-i\tau\varepsilon^{-1}A^{1/2}} - e^{-i\tau\varepsilon^{-1}(A^0)^{1/2}} \right) \varepsilon^{s} (H_0 + \varepsilon^{2}I)^{-s/2} \right\|_{L_2(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant C(\tau)\varepsilon.$$
(6)

Here $A = b(\mathbf{D})^* g(\mathbf{x}) b(\mathbf{D})$.

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Reduction 2: Direct integral expansion

Using the Floquet-Bloch theory, we expand the operators A, A^0 and H_0 (acting in $L_2(\mathbb{R}^d; \mathbb{C}^n)$) in the direct integrals:

$$A \sim \int_{\widetilde{\Omega}} \oplus A(\mathbf{k}) \, d\mathbf{k}, \quad A^0 \sim \int_{\widetilde{\Omega}} \oplus A^0(\mathbf{k}) \, d\mathbf{k}, \quad H_0 \sim \int_{\widetilde{\Omega}} \oplus H_0(\mathbf{k}) \, d\mathbf{k}.$$
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The parameter $\mathbf{k} \in \widetilde{\Omega}$ is called the *quasimomentum*. The operators $A(\mathbf{k})$, $A^{0}(\mathbf{k})$ and $H_{0}(\mathbf{k})$ act in $L_{2}(\Omega; \mathbb{C}^{n})$ and are defined by the expressions

$$\begin{aligned} A(\mathbf{k}) &= b(\mathbf{D} + \mathbf{k})^* g(\mathbf{x}) b(\mathbf{D} + \mathbf{k}), \\ A^0(\mathbf{k}) &= b(\mathbf{D} + \mathbf{k})^* g^0 b(\mathbf{D} + \mathbf{k}), \\ H_0(\mathbf{k}) &= |\mathbf{D} + \mathbf{k}|^2 \end{aligned}$$

with periodic boundary conditions. The precise definitions are given in terms of the corresponding quadratic forms.

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Reduction 2: Direct integral expansion

Using the direct integral expansions, we see that estimate

$$\left\| \left(e^{-i\tau\varepsilon^{-1}\mathcal{A}^{1/2}} - e^{-i\tau\varepsilon^{-1}(\mathcal{A}^{0})^{1/2}} \right) \varepsilon^{s} (\mathcal{H}_{0} + \varepsilon^{2}I)^{-s/2} \right\|_{L_{2}(\mathbb{R}^{d}) \to L_{2}(\mathbb{R}^{d})} \leqslant C(\tau)\varepsilon$$
(6)

is equivalent to estimate

$$\left\| \left(e^{-i\tau\varepsilon^{-1}\mathcal{A}(\mathbf{k})^{1/2}} - e^{-i\tau\varepsilon^{-1}\mathcal{A}^{0}(\mathbf{k})^{1/2}} \right) \varepsilon^{s} (\mathcal{H}_{0}(\mathbf{k}) + \varepsilon^{2}I)^{-s/2} \right\|_{L_{2}(\Omega) \to L_{2}(\Omega)} \leqslant \\ \leqslant C(\tau)\varepsilon, \quad (7)$$

for almost every $\mathbf{k} \in \widetilde{\Omega}$.

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Reduction 3: Projection onto the subspace of constant vector-valued functions

Next, let P be the projection onto the subspace

$$\mathfrak{N} = \{ \mathbf{u} \in L_2(\Omega; \mathbb{C}^n) : \mathbf{u} = \mathbf{c} \in \mathbb{C}^n \}.$$

Then

$$P \mathsf{u} = |\Omega|^{-1} \int_{\Omega} \mathsf{u}(\mathsf{x}) \, d\mathsf{x}.$$

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It is easily seen that for $s \ge 1$ estimate (7) is equivalent to

$$\left\| \left(e^{-i\tau\varepsilon^{-1}A(\mathbf{k})^{1/2}} - e^{-i\tau\varepsilon^{-1}S(\mathbf{k})^{1/2}} \right) P \right\|_{L_2(\Omega) \to L_2(\Omega)} \varepsilon^s (|\mathbf{k}|^2 + \varepsilon^2)^{-s/2} \leqslant C(\tau)\varepsilon^{-s/2}$$
(8)

for $|\mathbf{k}| \leq t^0$. Here

$$S(\mathbf{k}) = b(\mathbf{k})^* g^0 b(\mathbf{k}).$$

The operator $A(\mathbf{k})$ acting in $L_2(\Omega; \mathbb{C}^n)$ is an elliptic operator in a bounded domain; its spectrum is discrete. This operator depends on \mathbf{k} analytically. We consider A(0) as an unperturbed operator and $A(\mathbf{k})$ as a perturbed operator.

For $\mathbf{k} = 0$ the operator A(0) (given by $b(\mathbf{D})^*g(\mathbf{x})b(\mathbf{D})$ with periodic boundary conditions) has a kernel \mathfrak{N} consisting of constant vector-valued functions:

 $\operatorname{Ker} A(0) = \mathfrak{N} = \{ \mathbf{u} \in L_2(\Omega; \mathbb{C}^n) \colon \mathbf{u} = \mathbf{c} \in \mathbb{C}^n \}.$

So, the point $\lambda_0 = 0$ is an isolated eigenvalue of multiplicity *n* for A(0).

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So, the point $\lambda_0 = 0$ is an isolated eigenvalue of multiplicity *n* for A(0). Then for $|\mathbf{k}| \leq t^0$ the "perturbed" operator has exactly *n* eigenvalues $\lambda_1(\mathbf{k}), \ldots, \lambda_n(\mathbf{k})$ (counted with multiplicities) on $[0, \delta]$, and the interval $(\delta, 3\delta)$ is a gap in the spectrum of $A(\mathbf{k})$. (We control δ and t^0 .)



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Only these eigenvalues $\lambda_1(\mathbf{k}), \ldots, \lambda_n(\mathbf{k})$ (and the corresponding eigenfunctions) are important for our problem.

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Homogenization of hyperbolic equations

We put

$$\mathbf{k} = t\boldsymbol{\theta}, \quad t = |\mathbf{k}|, \quad \boldsymbol{\theta} \in \mathbb{S}^{d-1}.$$

and study the operator family $A(\mathbf{k}) = A(t\theta) =: A(t,\theta)$ by methods of the analytic perturbation theory with respect to the one-dimensional parameter t. But we have to make our constructions and estimates uniform in θ .

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and study the operator family $A(\mathbf{k}) = A(t\theta) =: A(t,\theta)$ by methods of the analytic perturbation theory with respect to the one-dimentional parameter t. But we have to make our constructions and estimates uniform in θ . By the Rellich–Kato theorem, for $t \leq t^0$ there exist real-analytic branches of eigenvalues $\lambda_l(t,\theta)$ and eigenvectors $\varphi_l(t,\theta)$ of the operator $A(t,\theta)$, l = 1, ..., n. We have

$A(t,\theta)\varphi_l(t,\theta) = \lambda_l(t,\theta)\varphi_l(t,\theta), \quad l = 1, \dots, n,$

and $\{\varphi_l(t, \theta)\}$ form an orthonormal basis in the eigenspace of $A(t, \theta)$ corresponding to $[0, \delta]$.

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Then we have the following power series expansions

$$\lambda_{l}(t,\theta) = \gamma_{l}(\theta)t^{2} + \mu_{l}(\theta)t^{3} + \dots, \quad l = 1,\dots,n,$$

$$\varphi_{l}(t,\theta) = \omega_{l}(\theta) + t\varphi_{l}^{(1)}(\theta) + \dots, \quad l = 1,\dots,n.$$

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Here $\gamma_l(\theta) \ge c_* > 0$ and $\mu_l(\theta) \in \mathbb{R}$. The "embryos" $\omega_1(\theta), \ldots, \omega_n(\theta)$ form an orthonormal basis in \mathfrak{N} .

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Here $\gamma_l(\theta) \ge c_* > 0$ and $\mu_l(\theta) \in \mathbb{R}$. The "embryos" $\omega_1(\theta), \ldots, \omega_n(\theta)$ form an orthonormal basis in \mathfrak{N} . However, we do not control the radius of convergence $t_*(\theta)$ of these expansions. The operator $S(\theta) = b(\theta)^* g^0 b(\theta)$ is called the *spectral germ* of the operator family $A(t, \theta)$ at t = 0.

Proposition 1 [M. Birman and T. Suslina, 2003]

 $S(\boldsymbol{\theta})\omega_l(\boldsymbol{\theta}) = \gamma_l(\boldsymbol{\theta})\omega_l(\boldsymbol{\theta}), \quad l = 1, \dots, n.$

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Threshold approximations

We need the so called threshold approximations for the spectral projection $F(t, \theta)$ of the operator $A(t, \theta)$ corresponding to $[0, \delta]$ and for the operator $A(t, \theta)F(t, \theta)$.

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Proposition 2 [M. Birman and T. Suslina, 2003]

Let $S(\theta) = b(\theta)^* g^0 b(\theta)$. For $t \leq t^0$ we have

$$\|F(t,\theta) - P\|_{L_2(\Omega) \to L_2(\Omega)} \leqslant C_1 t,$$
(9)
$$\|A(t,\theta)F(t,\theta) - t^2 S(\theta)P\|_{L_2(\Omega) \to L_2(\Omega)} \leqslant C_2 t^3$$
(10)

uniformly for $\theta \in \mathbb{S}^{d-1}$.

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uniformly for $\theta \in \mathbb{S}^{d-1}$.

Proposition 3 [M. Birman and T. Suslina, 2008]

 $\|A(t,\theta)^{1/2}F(t,\theta) - (t^2S(\theta))^{1/2}P\|_{L_2(\Omega) \to L_2(\Omega)} \leqslant C_3 t^2$ (11)

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We have

$$\left(e^{-i\tau A(t,\theta)^{1/2}} - e^{-i\tau(t^2 S(\theta))^{1/2}} \right) P = e^{-i\tau A(t,\theta)^{1/2}} (P - F(t,\theta)) P - (P - F(t,\theta)) e^{-i\tau(t^2 S(\theta))^{1/2}} P - - i \int_{0}^{\tau} e^{i(\tilde{\tau} - \tau)A(t,\theta)^{1/2}} \left(A(t,\theta)^{1/2} F(t,\theta) - (t^2 S(\theta))^{1/2} P \right) e^{-i\tilde{\tau}(t^2 S(\theta))^{1/2}} P \, d\tilde{\tau}.$$

Using estimates (9) and (11), we obtain

$$\left\|\left(e^{-i\tau A(t,\theta)^{1/2}}-e^{-i\tau(t^2S(\theta))^{1/2}}\right)P\right\|\leqslant 2C_1t+C_3|\tau|t^2,\quad t\leqslant t^0.$$

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In the estimate

$$\left\|\left(e^{-i\tau A(t,\theta)^{1/2}}-e^{-i\tau(t^2S(\theta))^{1/2}}\right)P\right\|\leqslant 2C_1t+C_3|\tau|t^2,\quad t\leqslant t^0.$$

we replace τ by $\tau \varepsilon^{-1}$ and multiply by the smoothing factor $\varepsilon^2(t^2 + \varepsilon^2)^{-1}$.

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we replace τ by $\tau \varepsilon^{-1}$ and multiply by the smoothing factor $\varepsilon^2 (t^2 + \varepsilon^2)^{-1}$. Then

$$\begin{split} \left\| \left(e^{-i\tau\varepsilon^{-1}A(t,\theta)} - e^{-i\tau\varepsilon^{-1}(t^2S(\theta))^{1/2}} \right) P \right\| \varepsilon^2 (t^2 + \varepsilon^2)^{-1} \leqslant \\ \leqslant (2C_1t + C_3\varepsilon^{-1}|\tau|t^2)\varepsilon^2 (t^2 + \varepsilon^2)^{-1} \leqslant (C_1 + C_3|\tau|)\varepsilon, \quad t \leqslant t^0. \end{split}$$

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This completes the proof. We arrive at the following result.

Theorem 1 [M. Birman and T. Suslina, 2008]

$$\|\cos(\tau A_{\varepsilon}^{1/2}) - \cos(\tau (A^0)^{1/2})\|_{H^2(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant (C_1 + C_3 |\tau|)\varepsilon.$$

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For more subtle considerations, we need more accurate approximation for $A(t, \theta)F(t, \theta)$:

Proposition 4 [M. Birman and T. Suslina, 2005]

For $t \leq t^0$ we have

$$A(t,\theta)F(t,\theta) = t^2 S(\theta)P + t^3 K(\theta) + \Psi(t,\theta).$$
(12)

The remainder term $\Psi(t, \theta)$ satisfies

$$\|\Psi(t,\theta)\|_{L_2(\Omega)\to L_2(\Omega)}\leqslant C_4t^4,\quad t\leqslant t^0.$$
(13)

The operator $K(\theta)$ is described in the invariant terms, as well as in terms of the coefficients of power series expansions for the eigenvalues and the eigenvectors of $A(t, \theta)$.

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For our problem, only the block of $K(\theta)$ in the subspace \mathfrak{N} is important. We have the following invariant representation:

$$\begin{split} \mathcal{N}(\boldsymbol{\theta}) &:= \mathcal{P}\mathcal{K}(\boldsymbol{\theta})\mathcal{P} = b(\boldsymbol{\theta})^* L(\boldsymbol{\theta}) b(\boldsymbol{\theta})\mathcal{P}, \\ L(\boldsymbol{\theta}) &:= |\Omega|^{-1} \int_{\Omega} \left(\Lambda(\mathbf{x})^* b(\boldsymbol{\theta})^* \widetilde{g}(\mathbf{x}) + \widetilde{g}(\mathbf{x})^* b(\boldsymbol{\theta}) \Lambda(\mathbf{x}) \right) \, d\mathbf{x}. \end{split}$$

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In terms of the coefficients,

$$\begin{split} \mathcal{N}(\theta) &= \mathcal{N}_{0}(\theta) + \mathcal{N}_{*}(\theta), \\ \mathcal{N}_{0}(\theta) &= \sum_{l=1}^{n} \mu_{l}(\theta)(\cdot, \omega_{l}(\theta))\omega_{l}(\theta), \\ \mathcal{N}_{*}(\theta) &= \sum_{l=1}^{n} \gamma_{l}(\theta)\left((\cdot, \widetilde{\omega}_{l}(\theta))\omega_{l}(\theta) + (\cdot, \omega_{l}(\theta))\widetilde{\omega}_{l}(\theta)\right), \end{split}$$

Here $\widetilde{\omega}_{l}(\boldsymbol{\theta}) \in \mathfrak{N}$ are defined by $\widetilde{\omega}_{l}(\boldsymbol{\theta}) := P\varphi_{l}^{(1)}(\boldsymbol{\theta})_{\mathbb{H}_{p}}$

Proposition 5 [M. Dorodnyi and T. Suslina, 2016]

For $t \leq t^0$ we have

 $A(t,\theta)^{1/2}F(t,\theta) = t S(\theta)^{1/2}P + t^2 G(\theta) + \Phi(t,\theta).$

The remainder term $\Phi(t, \theta)$ satisfies

 $\|\Phi(t,\theta)\|_{L_2(\Omega)\to L_2(\Omega)}\leqslant C_5t^3,\quad t\leqslant t^0.$

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The remainder term $\Phi(t, \theta)$ satisfies $\|\Phi(t, \theta)\|_{L_2(\Omega) \to L_2(\Omega)} \leq C_5 t^3, \quad t \leq t^0.$

Only the block of $G(\theta)$ in the subspace \mathfrak{N} is needed: $PG(\theta)P = \frac{1}{2}N_0(\theta)S(\theta)^{-1/2}P + S(\theta)^{-1/2}N_*(\theta)P +$

+ $N_*(\theta)S(\theta)^{-1/2}P + \mathcal{I}_*(\theta)P$,

where

$$\mathcal{I}_*(\theta) = -\frac{1}{\pi} \int_0^\infty s^{-1/2} (\Xi(t,s) N_*(\theta) + N_*(\theta) \Xi(t,s) - s \Xi(t,s) N_*(\theta) \Xi(t,s)) \, ds$$

and $\Xi(t,s) = (t^2 S(\theta) + s I)^{-1} P.$

Theorem 2 [M. Dorodnyi and T. Suslina, 2016]

Suppose that $N(\theta) = 0$ for any $\theta \in \mathbb{S}^{d-1}$. Then

 $\|\cos(\tau A_{\varepsilon}^{1/2}) - \cos(\tau (A^0)^{1/2})\|_{H^{3/2}(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant (\widetilde{C}_1 + \widetilde{C}_2|\tau|)\varepsilon.$ (14)

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The proof is similar to that of Theorem 1. We use more accurate threshold approximations.

Remark. If $A_{\varepsilon} = -\operatorname{div} g^{\varepsilon}(\mathbf{x})\nabla$, where $g(\mathbf{x})$ is a symmetric matrix with real entries, then $N(\theta) = 0$ for any $\theta \in \mathbb{S}^{d-1}$. Hence, (14) is true.

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Theorem 3 [M. Dorodnyi and T. Suslina, 2016]

Suppose that $N_0(\theta) = 0$ for any $\theta \in \mathbb{S}^{d-1}$ (this is equivalent to the assumption that $\mu_l(\theta) = 0$ for all l = 1, ..., n and any $\theta \in \mathbb{S}^{d-1}$). Suppose that the number p of different eigenvalues of $S(\theta)$ does not depend on θ . Then

$$\|\cos(\tau A_{\varepsilon}^{1/2}) - \cos(\tau (A^0)^{1/2})\|_{H^{3/2}(\mathbb{R}^d) \to L_2(\mathbb{R}^d)} \leqslant (\check{C}_1 + \check{C}_2|\tau|)\varepsilon.$$
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(15)

Corollary

Suppose that A_{ε} has real-valued coefficients. Suppose also that all eigenvalues of the germ $S(\theta)$ are simple for any $\theta \in \mathbb{S}^{d-1}$. Then estimate (15) is valid.

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Sharpness of [BSu, 2008] estimate in the general case

Finally, we confirm that Theorem 1 is sharp in the following sense.

Theorem 4 [M. Dorodnyi and T. Suslina, 2016]

Suppose that $N_0(\theta_0) \neq 0$ for some $\theta_0 \in \mathbb{S}^{d-1}$ (it means that $\mu_l(\theta_0) \neq 0$ for some *l*). Let $\tau \neq 0$ and s < 2. Then there does not exist a constant $C(\tau) > 0$ such that the estimate

 $\|\cos(au A_{arepsilon}^{1/2}) - \cos(au (A^0)^{1/2})\|_{H^s(\mathbb{R}^d) o L_2(\mathbb{R}^d)} \leqslant C(au) arepsilon$

holds for all sufficiently small $\varepsilon > 0$.

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holds for all sufficiently small $\varepsilon > 0$.

There are concrete examples of operators A_{ε} satisfying the assumptions of Theorem 4. One example is of the form $-\operatorname{div} g^{\varepsilon}(\mathbf{x})\nabla$, where $g(\mathbf{x})$ is Hermitian matrix with complex entries. Another example is the matrix operator with real-valued coefficients (the operator of elasticity theory in the cases of anisotropic and isotropic media, d = 2).

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Generalization. Applications

Similar results are obtained for more general operators of the form

 $\widetilde{A}_{\varepsilon} = (f^{\varepsilon}(\mathsf{x}))^* b(\mathsf{D})^* g^{\varepsilon}(\mathsf{x}) b(\mathsf{D}) f^{\varepsilon}(\mathsf{x}),$

where f(x) is a Γ -periodic $(n \times n)$ -matrix-valued function such that $f, f^{-1} \in L_{\infty}$.

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We apply the results to the following equations:

- The acoustics equation,
- The system of elasticity theory.

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